

Luminous Type II supernovae for their low expansion velocities

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ABSTRACT

We present optical and near-IR data of three Type II supernovae (SNe II), SN 2008bm, SN 2009aj, and SN 2009au. These SNe display the following common characteristics: signs of early interaction of the ejecta with circumstellar material (CSM), blue $B-V$ colours, weakness of metal lines, low expansion velocities, and V -band absolute magnitudes 2–3 mag brighter than those expected for normal SNe II based on their expansion velocities. Two more SNe reported in the literature (SN 1983K and LSQ13fn) share properties similar to our sample. Analysing this set of five SNe II, which are luminous for their low expansion velocities (LLEV), we find that their properties can be reproduced assuming ejecta-CSM interaction that lasts between 4–11 weeks post explosion. The contribution of this interaction to the radiation field seems to be the dominant component determining the observed weakness of metal lines in the spectra rather than the progenitor metallicity. Based on hydrodynamic simulations, we find that the interaction of the ejecta with a CSM of $\sim 3.6 M_{\odot}$ can reproduce the light curves and expansion velocities of SN 2009aj. Using data collected by the Chilean Automatic Supernova Search, we estimate an upper limit for the LLEV SNe II fraction to be 2–4 per cent of all normal SNe II. With the current data-set, it is not clear whether the LLEV events are a separated class of SNe II with a different progenitor system, or whether they are the extreme of a continuum mediated by CSM interaction with the rest of the normal SN II population.

Key words: circumstellar matter – supernovae: general – supernovae: individual: SN 1983K, SN 2008bm, SN 2009aj, SN 2009au, LSQ13fn

1 INTRODUCTION

Type II supernovae (SNe II; Minkowski 1941) are the final stage of the evolution of stars with an initial mass $> 8 M_{\odot}$, that retain a significant fraction of hydrogen in their envelopes at the moment of the collapse of their iron cores. Among SNe II, three special sub-types have been identified based on their photometric and spectroscopic characteristics: those showing hydrogen in early spectra that soon disappear (SNe IIb; Woosley et al. 1987; Filippenko 1988), those having light curves similar to SN 1987A (1987A-like SNe, e.g., Hamuy et al. 1988), and those showing narrow hydrogen emission lines in the spectra due to the interaction of the ejecta with a circumstellar material (CSM) (SNe IIn; Schlegel 1990). For the rest of SNe II (77 per cent; Shivvers et al. 2017, which show a range in luminosity decline rates) we will refer as normal SNe II.

Among SNe IIn, there are some events that after 25–50 d past explosion start to look similar to normal SNe II with broad P-Cygni profiles, but with bluer colours (hereafter SNe IIn/II, e.g., SN 1979C, Balinskaia et al. 1980, Branch et al. 1981, de Vaucouleurs et al. 1981; SN 1998S, Fassia et al. 2000, 2001; SN 2007pk, Inserra et al. 2013; SN 2008fq, Taddia et al. 2013, Faran et al. 2014; PTF11iqb, Smith et al. 2015; SN 2013fc, Kangas et al. 2016). On the other hand, Smith et al. (2015) suggested that some normal SNe II could be classified as SNe IIn if the classification spectra are taken at epochs early enough to detect the interaction of the ejecta with the CSM generated by progenitors with dense winds.

Recently, it has been shown that the presence of CSM around progenitors of normal SNe II seems to be ubiquitous (Förster et al. 2018). In fact, the ejecta-CSM interaction in normal SNe II, previously detected only in few cases (e.g., SN 2006bp; Quimby et al. 2007), now is confirmed in many cases by the detection of narrow hydrogen emission lines in very early-time spectroscopy (also called “flash spectroscopy”, Khazov et al. 2016; Yaron et al. 2017). The effect of this early ejecta-CSM interaction over the radiation field is also inferred from the light curve modelling (Moriya et al. 2011; Morozova et al. 2017, 2018; Das & Ray 2017; Förster et al. 2018) and, more recently, from the simultaneous modelling of light curves and spectra (Hillier & Dessart 2019).

With the increasing number of discovered SNe II, more rare events are being revealed. An example of this is LSQ13fn (Polshaw et al. 2016), an SN II that shares some of the properties seen on SNe IIn/II (e.g., early ejecta-CSM interaction, blue colours, and high luminosities) but with two additional characteristics: (1) a weakness of metal lines in the spectra, and (2) low expansion velocities, with values between the subluminous SN II 2005cs (Pastorello et al. 2006, 2009) and the archetypal SN II 1999em (Hamuy et al. 2001; Leonard et al. 2002; Elmhamdi et al. 2003), but having a luminosity comparable to the moderately luminous SN II 2004et (Sahu et al. 2006; Maguire et al. 2010). In addition to LSQ13fn, SN 1983K was also reported showing some of the aforementioned characteristics (Niemela et al. 1985; Phillips et al. 1990). SN 1983K and LSQ13fn, being more luminous than the expected from their expansion velocities, do not fall on the luminosity-expansion velocity relation found by Hamuy & Pinto (2002) for normal SNe II.

Polshaw et al. (2016) suggested a combined effect of a residual thermal energy from the early ejecta-CSM interaction, a low metallicity, and a large radius of the progenitor to explain the peculiarities seen in LSQ13fn. In this work we present optical and near-IR data of SN 2008bm, SN 2009aj, and SN 2009au, which show similar properties to those seen in SN 1983K and LSQ13fn. Throughout this paper we identify this sample of SNe II using the acronym Luminous for their Low Expansion Velocities (LLEV) SNe II.

The paper is organized as follows. In Section 2 we present the relevant information on our SN data-set and their host galaxies. Photometry along with the description of the data reduction is presented in Section 3. In Section 4 and 5, we contrast the properties of the LLEV SNe II with those of other SNe II. A possible scenario that could explain the peculiar characteristics seen in the LLEV SNe II is discussed in Section 6. We also present an estimation of the fraction of LLEV SNe II with respect to normal SNe II in Section 6.3 and their impact on the use of normal SNe II as distance indicators in Section 6.4. Finally, our conclusions are summarized in Section 7.

2 SUPERNOVAE AND HOST GALAXIES

In Table 1 we list the main parameters of the LLEV SNe II in our set and their host galaxies. Throughout this work we assume the extinction curve given by Fitzpatrick (1999) with $R_V = 3.1$. To compute distances from recession redshifts, we assume $H_0 = 73 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_m = 0.27$, and $\Omega_{\Lambda} = 0.73$. To estimate the error in distances due to peculiar velocities, we include a velocity dispersion of 382 km s^{-1} (Wang et al. 2006).

2.1 SN 2008bm

SN 2008bm was discovered in the galaxy CGCG 71–101 on 2008 March 29.3 UT (Drake et al. 2008) by the Catalina Real-time Transient Survey (CRTS; Drake et al. 2009). The SN is also present on images taken on January 31.5 UT. The last non-detection was on January 11.5 UT, therefore the explosion epoch is constrained to occur at MJD 54486.5 ± 10.0^1 , which is 67.8 d before the discovery. The event was classified as an SN IIn a couple of months past explosion based on a spectrum obtained on April 7.1 UT. Stritzinger & Morrell (2008) initially classified SN 2008bm as an SN IIn, while latter spectra reported by Gutiérrez et al. (2017a, hereafter G17) show clear Balmer absorption lines, typically seen in normal SNe II. No radio detection was obtained in the 8.46 GHz band at the SN position on December 07.7 UT (Chandra & Soderberg 2008).

CGCG 71–101 has a recession velocity of 9875 km s^{-1} (NED²), which translates into a distance modulus of $35.66 \pm 0.08 \text{ mag}$. No NaID absorption at the redshift of the host galaxy was detected in the SN spectra, indicating a negligible host galaxy reddening ($E_{B-V}^{\text{host}} = 0.00 \pm 0.03 \text{ mag}$,

¹ We assume the explosion epoch as the midpoint of the range between the last non-detection and the first detection of the SN, with the error (not normal but uniform) being half the range.

² NASA/IPAC Extragalactic Database (<https://ned.ipac.caltech.edu/>).

Table 1. SN and host galaxy parameters of the LLEV SN II sample.

SN data	1983K	2008bm	2009aj	2009au	LSQ13fn
RA (J2000.0)	12:46:36.41	13:02:58.78	13:56:45.33	12:59:46.00	11:51:17.29
DEC (J2000.0)	−8:21:21.9	+10:30:27.0	−48:29:36.2	−29:36:07.5	−29:36:41.10
Host galaxy	NGC 4699	CGCG 71–101	ESO 221–G18	ESO 443–G21	LEDA 727284
Host type ^a	SAB(rs)b	Sc	Sa?	Scd	Sa ^e
Distance modulus (mag)	31.50 ± 0.35	35.66 ± 0.08	33.89 ± 0.20	33.31 ± 0.22	37.21 ± 0.05
E_{B-V}^{MW} (mag) ^b	0.015 ± 0.002	0.022 ± 0.003	0.124 ± 0.020	0.079 ± 0.013	0.054 ± 0.009
E_{B-V}^{host} (mag)	0.0 ^d	0.00 ± 0.03	0.00 ± 0.02	0.35 ± 0.17	0.0 ^e
Explosion epoch (MJD)	45490.1 ± 1.0	54486.5 ± 10.0	54879.8 ± 6.5	54897.2 ± 4.0	56299.2 ± 1.0
SN heliocentric velocity (km s ^{−1})	1394 ± 162	9628 ± 25	2883 ± 162	2875 ± 30	18900 ± 300 ^e
12+log(O/H) (dex) ^c	8.55 [†]	8.33 ± 0.02 [*]	8.29 ± 0.03 [*]	8.56 ± 0.03 [*]	< 8.50 ± 0.16 [⊗]
s_2 (mag (100 d) ^{−1})	0.73 ± 0.10	2.56 ± 0.21	0.80 ± 0.02	3.06 ± 0.02	0.86 ± 0.10
M_V^{max} (mag)	−18.83 ± 0.36	−18.40 ± 0.18	−18.81 ± 0.21	−17.55 ± 0.58	−17.95 ± 0.21
M_V^{50d} (mag)	−18.43 ± 0.35	−18.33 ± 0.15	−17.81 ± 0.22	−16.54 ± 0.59	−17.33 ± 0.07
⁵⁶ Ni mass (M _⊙)	0.056 ± 0.018	> 0.015 ± 0.005	> 0.043 ± 0.010	–	> 0.018 ± 0.005
(B−V) _{50d} (mag)	0.30 ± 0.03	< 0.12 ± 0.03	0.41 ± 0.05	0.48 ± 0.15	0.34 ± 0.11
pEW(Fe II λ5018) _{50d} (Å)	3.46 ± 0.22	< 8.00 ± 3.00	9.02 ± 0.90	14.65 ± 2.63	9.00 ± 2.20
v_{FeII}^{50d} (km s ^{−1})	2331 ± 189	1684 ± 110	2336 ± 181	1549 ± 40	2320 ± 345

^a From NED, unless otherwise noted.

^b Galactic colour excesses from [Schlafly & Finkbeiner \(2011\)](#), with a statistical error of 16 per cent ([Schlegel et al. 1998](#)).

^c Oxygen abundances, in the N2 calibration of [Marino et al. \(2013\)](#), measured by [Kuncarayakti et al. \(2018\)](#) (†), [Anderson et al. \(2016\)](#) (★), or [Polshaw et al. \(2016\)](#) (⊗). The latter is a recalibration of the original value reported in the N2 calibration of [Pettini & Pagel \(2004\)](#).

^d Value from [Niemela et al. \(1985\)](#).

^e Value from [Polshaw et al. \(2016\)](#).

[Anderson et al. 2014b](#), hereafter A14). SN 2008bm is located at a projected distance of 9.3 ± 0.3 kpc from the center of the apparently nearly face-on host galaxy, which is consistent with the low reddening scenario.

2.2 SN 2009aj

SN 2009aj was discovered in the galaxy ESO 221–G18 on 2009 February 24.3 UT ([Pignata et al. 2009b](#)) during the Chilean Automatic Supernova Search (CHASE; [Pignata et al. 2009a](#)). Nothing was visible at the SN position on February 11.2 UT, therefore the explosion is constrained to occur at MJD 54879.8 ± 6.5 (i.e., 6.5 d before the discovery). The event was classified by [Stritzinger et al. \(2009\)](#) as an SN II around maximum and reminiscent of SN 1983K.

We do not detect the presence of Na I D at the redshift of the host galaxy in the SN spectra. In this case, we assume zero reddening with an error corresponding to the 3σ upper limit of the Na I D pseudo equivalent width (pEW) non-detection. Using the relation of [Poznanski et al. \(2012\)](#)³ we obtain $E_{B-V}^{host} = 0.00 \pm 0.02$ mag, which we adopt as the host galaxy colour excess. ESO 221–G18 has a recessional velocity corrected for the infall of the Local Group toward the Virgo cluster and the Great Attractor of 4380 ± 112 km s^{−1} (NED), which translates into a distance modulus of 33.89 ± 0.20 mag.

2.3 SN 2009au

SN 2009au was discovered in the galaxy ESO 443–G21 on 2009 March 11.2 UT ([Pignata et al. 2009c](#)) by the CHASE survey. Nothing was visible at the SN position on March

3.2 UT, therefore the explosion epoch is constrained to occur at MJD 54897.2 ± 4.0 (i.e., 4 d before the discovery). The event appeared to be a young SN IIn soon after explosion ([Stritzinger et al. 2009](#)), however later spectra (see G17) show clear Balmer absorption features. No radio detection was obtained at the SN position in the 8.46 GHz band on December 13.6 UT ([Chandra & Soderberg 2009](#)).

The distance modulus to ESO 443–G21 is estimated using the Tully-Fisher relation to be 33.44 ± 0.45 mag ([Tully et al. 2016](#)). In addition, ESO 443–G21 is a member of the galaxy group HDCE 754 ([Crook et al. 2007](#)), which has a recessional velocity corrected for Virgo infall of 3287 km s^{−1} (NED), corresponding to a distance modulus of 33.27 ± 0.25 mag. We adopt the weighted mean of these values ($\mu = 33.31 \pm 0.22$ mag) as the distance to ESO 443–G21.

SN 2009au is located at a projected distance of 1.7 ± 0.1 kpc from the center of its edge-on host galaxy, so the SN could be affected by a high amount of extinction. We measured a Na I D pEW of 1.33 ± 0.21 Å at the redshift of the host galaxy in the SN spectra. However, the Na I D pEW becomes insensitive to estimate reddening for pEW > 1.0 Å ([Phillips et al. 2013](#)). Despite the above, we can estimate a lower limit for the E_{B-V}^{host} value as the saturation point (pEW = 1.0 Å) in the [Poznanski et al. \(2012\)](#) relation ($E_{B-V}^{host} > 0.21 \pm 0.14$ mag). On the other hand, matching the B−V colour curve of SN 2009au to the rest of LLEV SNe II we obtain $E_{B-V}^{host} = 0.35 \pm 0.17$ mag, which is consistent with the previous lower limit. We adopt the colour excess from the B−V colour match as the E_{B-V}^{host} value for SN 2009au.

³ Where, as noted by [Phillips et al. \(2013\)](#), the error around the log(E_{B-V}) relation is of 0.30 dex.

3 OBSERVATIONAL MATERIAL

3.1 Photometric data

Optical and near-IR images of SN 2008bm, SN 2009aj, and SN 2009au were obtained over the course of the *Carnegie Supernova Project I* (CSP-I; Hamuy et al. 2006). Johnson *BV* and Sloan *ugri* images were obtained mostly with the 1 m Swope telescope at Las Campanas Observatory (LCO), while near-IR *YJH* images were obtained with both the 1 m Swope and the LCO 2.5 m du Pont telescope. The CSP-I data reduction is described in Contreras et al. (2010) and Krisciunas et al. (2017).

In addition, Johnson-Kron-Cousins *UBVRI* and Sloan *griz* images of SN 2009aj and SN 2009au were also obtained with the 41 cm Panchromatic Robotic Optical Monitoring and Polarimetry Telescopes (PROMPT; Reichart et al. 2005) at Cerro Tololo Inter-American Observatory (CTIO), as part of the CHASE follow-up program. We also include *BVRI* images of SN 2009au obtained with the SMARTS 1.3 m telescope at CTIO, equipped with the A Novel Dual Imaging Camera (ANDICAM)⁴. Data reduction includes bias subtraction, overscan correction, flat-field correction, cosmic ray rejection, and image combination.

As part of the CSP-I and CHASE programs, images of the SNe host galaxies 1.5–4.0 years after the SN explosion were obtained in order to remove the host galaxy contamination. In the case of the ANDICAM images of SN 2009au, to construct a template for each filter, we selected the earliest images of the SN field obtained under good observing conditions with such camera. Then, using the SNOOPY⁵ package, we removed the SN flux. Such templates were then subtracted from the rest of the images acquired with ANDICAM. Instrumental magnitudes were measured using the point spread function (PSF) technique. Final magnitudes were computed with respect to a local sequence of stars (Tables A1–A6), which are calibrated using Landolt (1992) *UBVRI*, Smith et al. (2002) *ugriz*, and Persson et al. (1998) *JHK* standard stars (*J*- and *K*-band observations were used to derive *Y*-band magnitudes of the standard stars, see Hamuy et al. 2006).

Tables A7–A9 list the standard system photometry of SN 2008bm, SN 2009aj, and SN 2009au, while Fig. 1 shows their optical and near-IR light curves. Systematic differences between the CSP-I and CHASE *BVgr* photometry, CSP-I and ANDICAM *BV* photometry, and CHASE and ANDICAM *RI* photometry are, on average, lower than 0.02 mag. We stress the fact that the photometry obtained with the SMARTS 1.3 m and PROMPT telescopes was calibrated using the magnitudes of the CSP-I local sequences.

In the analysis we include the unfiltered photometry transformed to *V*-band magnitudes (V_{CRTS}) of SN 2008bm obtained by the CRTS⁶. We note that the V_{CRTS} photometry lies very close to our *r*-band photometry in the phase interval where the light curves overlap. In order to reach a better

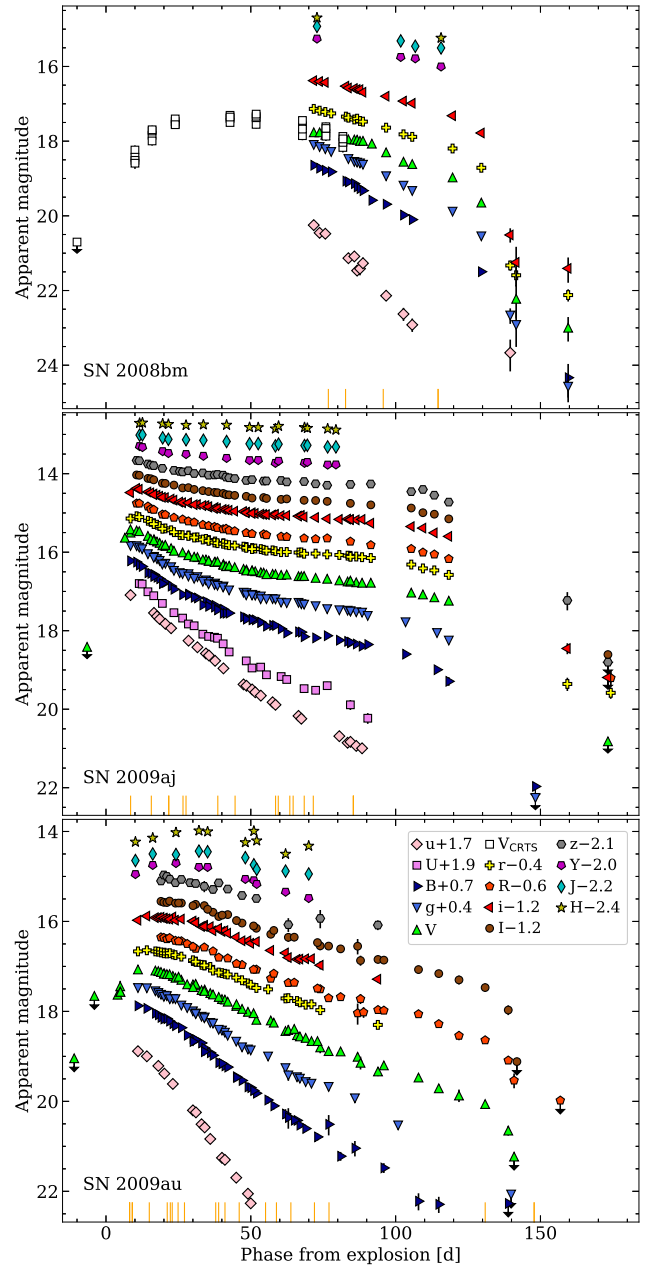


Figure 1. Light curves of SN 2008bm (top), SN 2009aj (middle), and SN 2009au (bottom). Bands and magnitude shifts with respect to the original values are quoted in the legend. Arrows indicate upper limits. Orange ticks mark the epochs of the spectroscopy. For SN 2008bm we also plot the CRTS photometry.

agreement between the V_{CRTS} and the filtered *V*-band photometry, we add 0.26 mag to the V_{CRTS} magnitudes, which is the average $V-r$ colour at 50 d since explosion we obtained for SN 2009aj and LSQ13fn.

3.2 Spectroscopic data

Optical spectra of SN 2009au were obtained with telescopes and instruments listed in Table A10. Data reduction includes bias subtraction, flat-field correction, wavelength calibration, one-dimensional spectrum extraction and sky sub-

⁴ Images are available on the NOAO science archive (<http://archive1.dm.noao.edu/>).

⁵ SNOOPY is a package for SN photometry using PSF fitting and/or template subtraction developed by E. Cappellaro. A package description can be found at <http://sngroup.oapd.inaf.it/snoopy.html>

⁶ Available on <http://nessi.cacr.caltech.edu/DataRelease/CRTS-I> or <http://archive.noao.edu>

traction, and flux calibration. The left part of Fig. A1 shows the optical spectral evolution of SN 2009au.

We also obtained near-IR spectra of SN 2009aj with telescopes and instruments listed in Table A11. Data reduction includes the subtraction of the pairs of images to remove the sky background, images combination, wavelength calibration, one-dimensional spectrum extraction, telluric correction, and flux calibration. The right part of Fig. A1 shows the near-IR spectral evolution of SN 2009aj.

In the analysis we include the optical spectroscopy for SN 2008bm, SN 2009aj, and SN 2009au, which was obtained by the CSP-I and already published in G17.

3.3 Sample of supernovae

In addition to SN 2008bm, SN 2009aj, and SN 2009au, we include SN 1983K (Niemela et al. 1985; Phillips et al. 1990) and LSQ13fn (Polshaw et al. 2016) into the analysis, given the similarity of their photometric and spectral properties (see Section 4 and 5). SN 1983K was discovered in NGC 4699 on 1983 June 6.1 UT (Maza et al. 1983) at 17.1 mag (Phillips et al. 1990). There is no information about the last non-detection. However, the SN was at 13.3 mag on June 10.1 UT (Phillips et al. 1990), which indicates that the SN was discovered close to the explosion. Fitting a quadratic polynomial to the B -band rise photometry, we obtain the explosion epoch to be MJD 45490.1 \pm 1.0 (i.e., 1 d before the discovery). The distance modulus to NGC 4699 is estimated with the Tully-Fisher relation to be 31.45 \pm 0.45 mag (Tully et al. 2016). In addition, NGC 4699 is a member of the galaxy group HDCE 740 (Crook et al. 2007), which has a recessional velocity corrected for Virgo infall of 1506 km s⁻¹ (NED), corresponding to a distance modulus of 31.57 \pm 0.55 mag. We adopt the weighted mean ($\mu = 31.50 \pm 0.35$ mag) as the distance to NGC 4699. LSQ13fn was discovered in LEDA 727284 on 2013 January 10.2 UT by the La Silla-Quest Variability Survey (Baltay et al. 2013), which also obtained early-time photometry. The first detection was on January 8.2 UT, while the last non detection was on January 6.2 UT, so the explosion is constrained to occur at MJD 56299.2 \pm 1.0 (i.e., 3 d before the discovery). Using the redshift given in Polshaw et al. (2016) and the cosmic microwave background dipole model of Fixsen et al. (1996), we compute a recessional velocity of 19225 \pm 300 km s⁻¹ for LEDA 727284, corresponding to a distance modulus of 37.21 \pm 0.05 mag. The main parameters for SN 1983K, LSQ13fn, and their host galaxies are summarized in Table 1.

4 PHOTOMETRIC PROPERTIES

In this section we compare photometric properties of the LLEV SNe II with those of other SNe II.

4.1 V -band light curves

Fig. 2a shows the absolute V -band light curves of the LLEV SNe II in our set. We note that the length of the photospheric phase (also called optically thick phase duration, OPTd) of SN 2008bm and SN 2009aj is about 120 d, and of about 130 d for SN 2009au. These values are comparable to that

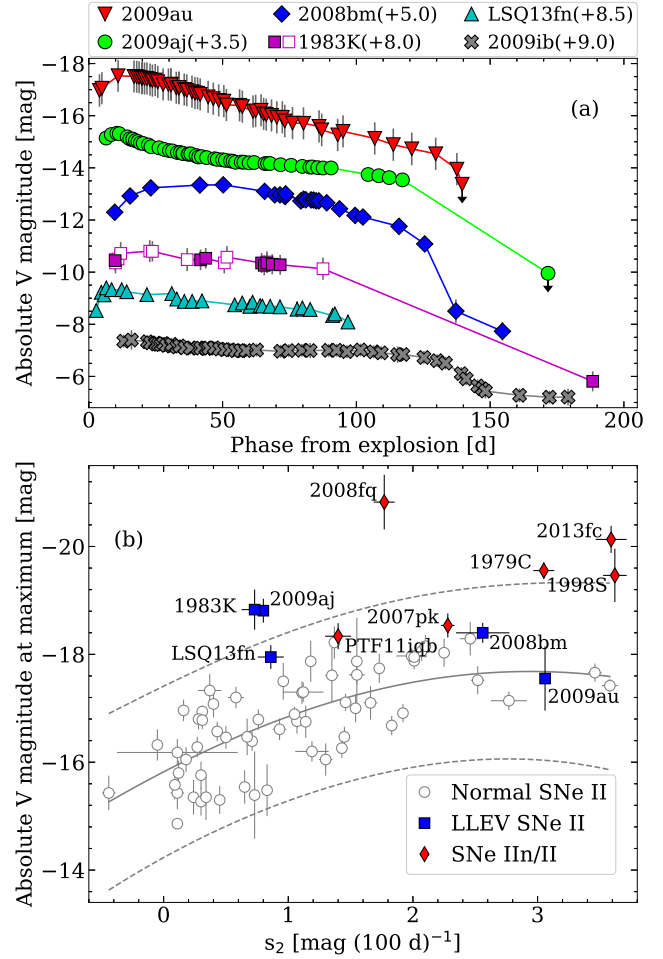


Figure 2. (a) Absolute V -band light curves of the LLEV SNe II in our set (coloured symbols), and of the normal SN II 2009ib (gray crosses). For clarity we averaged photometric points into 2 hr bins, and applied and offset to magnitudes (indicated in parentheses). Arrows indicate upper limits. (b) Absolute V -band magnitude at maximum versus V -band decline rate during the plateau, showing the LLEV SNe II in our set (blue squares), the normal SNe II in the A14 sample corrected for E_{B-V}^{host} (empty circles), and the SNe II in/II from the literature (red diamonds). The solid line corresponds to the Gaussian process fit, where dashed lines indicate the 3 σ error around the fit.

of SNe II 2004er (120 d; Anderson et al. 2014b) and 2009ib (Takáts et al. 2015, gray crosses), which are among the normal SNe II with longest plateau.

Analysing a set of 116 SN II V -band light curves, A14 found a correlation between the absolute V -band magnitude at maximum (M_V^{max}) and the decline rate of the second, shallower slope in the light curve (s_2), which suggests a continuum in the normal SN II population in this parameter space. Similar peak magnitude-decline rate correlations were also obtained by Sanders et al. (2015), Galbany et al. (2016), and Valenti et al. (2016).

For SN 2008bm, SN 2009aj, SN 2009au, and LSQ13fn we have V -band photometry around the maximum light. Therefore we can check if such correlation holds also for LLEV SNe II. In the case of SN 1983K, the V -band light curve (purple filled squares) is less sampled than the B -band one (see Fig. 3 of Phillips et al. 1990), and the maxi-

mum light (observed in B -band) in the V -band is missed. To estimate the V -band photometry around the maximum, we interpolate the $B-V$ colour (which is monotonically increasing, see Section 4.3) to the epochs of the B -band photometry without V -band measurements (t^*), and then we compute $V_{t^*} = B_{t^*} - (B-V)_{t^*}$ (purple empty squares). To estimate the V -band maximum of our LLEV SNe II, we fit a fourth order polynomial to the photometry close in time to the brightest point. To estimate s_2 , we fit a straight line to the V -band light curve during the plateau phase (for more details, see A14). Values of M_V^{\max} and s_2 for the LLEV SNe II set are listed in Table 1.

Fig. 2b shows the normal SNe II in the A14 sample⁷ (empty circles; where values are from Table 6 of A14) in the M_V^{\max} versus s_2 space, where we remove those SNe without estimation of E_{B-V}^{host} . To characterize the distribution of the A14 sample in this space, we perform a Gaussian process fit (solid line), where the dashed lines indicate the $\pm 3\sigma$ limits. In the figure we plot the LLEV SNe II as blue squares. We see that they have $M_V^{\max} \lesssim -17.5$ mag, where SN 2008bm and SN 2009au are fast decliners, while for SN 1983K, SN 2009aj, and LSQ13fn the decline rate is lower, indicating that LLEV SNe II seems not to be connected to a particular light curve decline rate. In addition, we see that SN 2008bm, SN 2009au, and LSQ13fn are within 3σ limit, while SN 1983K and SN 2009aj are outliers in the distribution, i.e., they are significantly brighter than what would be implied from their s_2 decline rates. In the figure we also plot the SNe IIn/II (red diamonds) that we found in the literature (the set and main properties are listed in Table A12). We see that the LLEV SNe II tend to have lower decline rates than SNe IIn/II.

The case of SN 2009au is special. Its OPTd and s_2 values are not common in normal SNe II, where faster decliners have also shorter OPTd (see for example, Fig.12 in A14). Performing radiation hydrodynamics simulations, Hillier & Dessart (2019) showed that the ejecta-CSM interaction scenario could produce fast decliners with an OPTd similar to slow decliners. Based on the latter, we suggest that the characteristics of the V -band light curve of SN 2009au are consequence of the interaction of its ejecta with a CSM.

4.2 Pseudo-bolometric light curves

To compute pseudo-bolometric light curves, we proceed as follows: (1) We convert broad-band magnitudes into monochromatic fluxes (f_x), associated to their respective effective wavelengths (l_x). For a fair comparison with other SNe II in the literature, we use $BVRI$ and gri magnitudes. (2) For each photometric epoch, we perform linear interpolations between the (l_x , f_x) points, which we adopt as the spectral energy distribution (SED). (3) We correct the SED for redshift and colour excess, and then we integrate it from 4200 to 7500 Å (which is a wavelength range covered by the filters). (4) We convert the integrated flux into luminosity using the corresponding SN distance.

The top part of Fig. 3 shows the $BVRI$ pseudo-bolometric light curves of the LLEV SNe II in our sam-

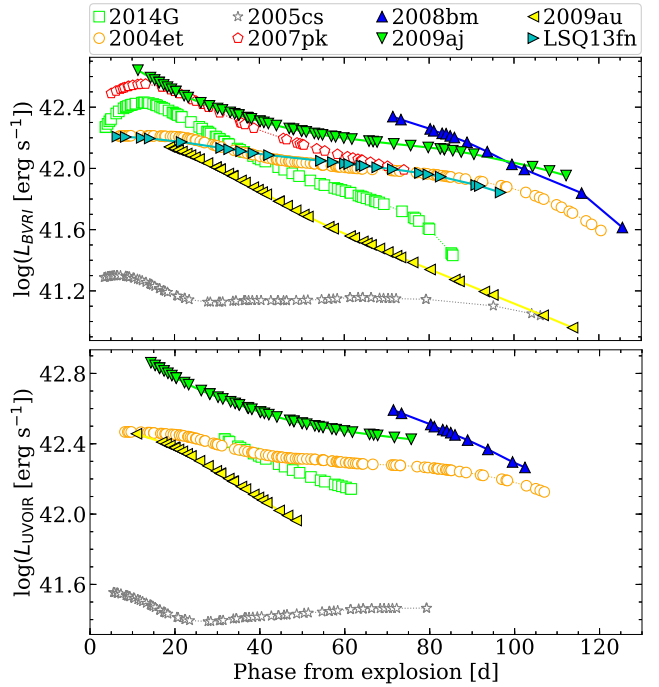


Figure 3. $BVRI$ (top) and $UVOIR$ (bottom) pseudo-bolometric light curves of the LLEV SNe II in our set (filled symbols) and a selection of SNe II that share some of the properties seen in LLEV SNe II (empty symbols).

ple, except SN 1983K for which the photometry in the necessary bands is not available. For comparison, we also compute the $BVRI$ pseudo-bolometric light curves for a selection of SNe II from the literature that share some of the properties observed in our LLEV SN II sample: normal SNe II 2014G (Bose et al. 2016; Terreran et al. 2016) and SN 2004et, which have decline rates similar to SN 2008bm and SN 2009au, and to SN 2009aj and LSQ13fn, respectively; the SN IIn/II 2007pk; and the subluminal SN II 2005cs, which has low expansion velocities (see Section 5.3). We see that SN 2008bm and SN 2009aj have luminosities higher than SNe II with a similar decline rate, confirming the result obtained in Fig. 2b, where the luminosity of SN 2009aj is comparable to those of the SN 2007pk at epochs earlier than 40 d since explosion. After that, the evolution of the SN 2009aj luminosity is similar to that of SN 2004et, though brighter. SN 2009au is less luminous than SN 2014G in all the phase range reported in the plot, with a luminosity at maximum light similar to SN 2004et, and similar to SN 2005cs at 100 d. LSQ13fn, as reported by Polshaw et al. (2016), has a luminosity similar to SN 2004et. We also compute the pseudo-bolometric light curve using ultraviolet, optical, and near-IR ($UVOIR$) photometry, from u/U - up to H -band (3600–16500 Å, bottom part of Fig. 3). We find that SN 2008bm and SN 2009aj are not only brighter in the optical (top part of Fig. 3), but also in the $UVOIR$ wavelength range.

4.3 Colour evolution

Normal SNe II are found to form a continuum population in observed colours, showing a large diversity at all epochs.

⁷ We remove SN 2008bm and SN 2009au from the A14 sample since they are in our LLEV SN II sample.

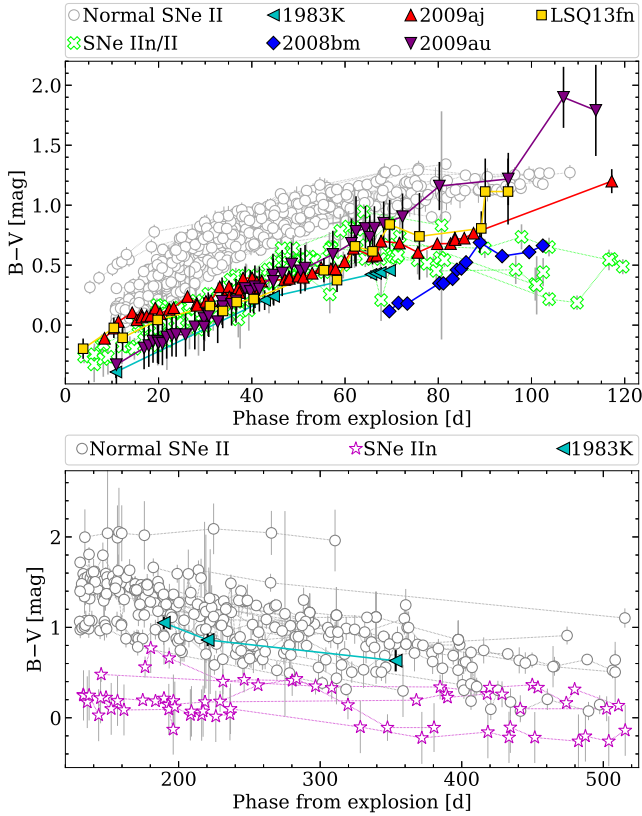


Figure 4. Evolution of the $B-V$ colour, corrected for E_{B-V}^{host} . Top: during the first 120 d since explosion of the LLEV SNe II in our set (filled symbols) compared with the normal SNe II in the D18 sample (empty circles). Bottom: during the late-time phase of SN 1983K (filled triangle) and of normal SNe II from the literature with B - and V -band photometry during 190–350 d since explosion (empty circles). We also include SNe IIn/II (top, green crosses) and SNe IIn (bottom, magenta stars) to depict the effect of the ejecta-CSM interaction over the $B-V$ colour.

While there are some (red) normal SNe II that clearly show the effects of strong host galaxy reddening, most of the colour dispersion apparently arises from intrinsic colour differences between normal SNe II (de Jaeger et al. 2018, hereafter D18).

The top part of Fig. 4 shows the $B-V$ colour curves of the normal SNe II analysed in D18⁸ (empty circles) and corrected for E_{B-V}^{host} (using values of A14), along with the LLEV SNe II in our set. We note that the LLEV SNe II are systematically bluer than normal SNe II in all the phase range reported in the plot⁹. The latter could be due to low metallicity progenitors (where the lower the progenitor metallicity, the bluer the colour; Dessart et al. 2014), or a consequence of an ejecta-CSM interaction, which makes colours bluer than normal SNe II (e.g., Hillier & Dessart 2019) as in the case of SNe IIn/II, which are included in Fig. 4 for comparison.

The bottom part of Fig. 4 shows the evolution of the

late-time $B-V$ colour of SN 1983K, compared with normal SNe II that we found in the literature with B - and V -band photometry during 190–350 d since explosion, i.e., the phase range covered by the late-time $B-V$ colour of SN 1983K. This sample is listed in Table A13. We see that the $B-V$ colour evolution of SN 1983K is consistent with the rest of normal SNe II in the plot. We also include a set of SNe IIn (sample listed in Table A14) to depict the effect of a long lasting ejecta-CSM interaction over the $B-V$ colours. We see that the ejecta-CSM interaction on SNe IIn makes their colour bluer than normal SNe II. The latter suggests that for SN 1983K, the late-time flux is dominated by the decay of ^{56}Co to ^{56}Fe .

4.4 Nickel mass

Assuming that the observed flux during the radioactive tail is due to the decay of ^{56}Co to ^{56}Fe , and that all the γ -rays from that decay are thermalized, the ^{56}Ni mass synthesized during a SN II explosion can be estimated as

$$\frac{M(^{56}\text{Ni})}{M_{\odot}} = \left(\frac{L_t}{L_*} \right) \exp \left[\frac{(t - t_0)/(1 + z) - 6.1 \text{ d}}{111.26 \text{ d}} \right] \quad (1)$$

(e.g., Hamuy 2003), where L_t is the bolometric luminosity measured at epoch t , $L_* = 1.271 \times 10^{43} \text{ erg s}^{-1}$, and t_0 and z are the explosion epoch and heliocentric redshift of the SN, respectively.

To estimate L_t for the LLEV SNe II in our set, we use

$$\log(L_t/L_*) = (\mu - V_{0,t} - \text{BC})/2.5 - 7.38, \quad (2)$$

where $V_{0,t}$ is the V -band magnitude at epoch t during the radioactive tail corrected by extinction, μ is the SN distance modulus, and BC is the bolometric correction. For the latter, we assume the same BC for SNe II during the radioactive tail ($\text{BC} = 0.26 \pm 0.06 \text{ mag}$; Hamuy 2001). For SN 2009aj we convert the last Rr photometry (at 173 d since explosion) to V -band magnitude using $V-R$ and $V-r$ colours of normal SNe II nearly at the same epoch of the Rr photometry, obtaining $V = 20.60 \pm 0.13 \text{ mag}$. For LSQ13fn, we compute $V = 26.24 \pm 0.30 \text{ mag}$ from its nebular spectrum (at 335 d since explosion). For SN 1983K we measure a V -band radioactive tail slope (s_3) of $1.01 \pm 0.02 \text{ mag (100 d)}^{-1}$, which is consistent with the slope of $0.98 \text{ mag (100 d)}^{-1}$ expected for the complete γ -ray trapping scenario. For SN 2008bm, SN 2009aj, and LSQ13fn, there are not enough photometric data during the nebular phase to estimate s_3 . Since we cannot check the complete γ -ray trapping scenario for those SNe, we adopt their ^{56}Ni mass estimations as lower limits. ^{56}Ni mass values are listed in Table 1

Fig. 5 shows the location of the LLEV SNe II (blue squares) in the $\log(M(^{56}\text{Ni}))$ versus the absolute V -band magnitude at 50 d since explosion ($M_V^{50\text{d}}$) space¹⁰. For comparison, we include normal SNe II (empty circles) in the A14 sample, which have ^{56}Ni mass values estimated in the same manner than our LLEV SNe II, and the normal SNe II in the Hamuy (2003) sample, where we recompute the ^{56}Ni masses using new estimations for the distance

⁸ We remove SN 2009au from the D18 sample since it is in our LLEV SN II sample.

⁹ In the case of SN 2009au, its $B-V$ colour curve was matched to the rest of LLEV SNe II in order to estimate E_{B-V}^{host} (see Section 2.3).

¹⁰ Since the errors in the $\log(M(^{56}\text{Ni}))$ versus $M_V^{50\text{d}}$ space are dominated by errors in reddenings and/or distances, the confidence regions are elongated ellipsoids (e.g., Pejcha & Prieto 2015).

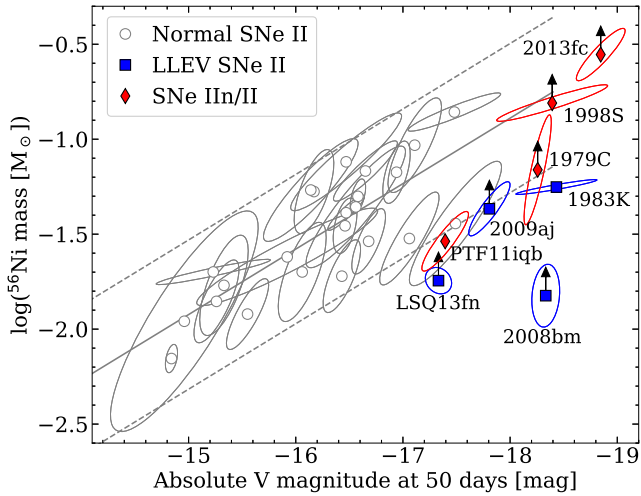


Figure 5. ^{56}Ni mass versus absolute V -band magnitude at 50 d since explosion, showing the normal SNe II (empty circles), the SNe IIn/II from the literature (red diamonds), and the LLEV SNe II in our set (blue squares). The ellipses indicate the 1σ statistical error. The solid line corresponds to a linear fit, where dashed lines indicate the 2σ error around the fit. Arrows indicate lower limits.

and host galaxy reddening (the SN set and parameters are listed in Table A15). We also include the SN IIn/II set (red diamonds). As noted by Hamuy (2003) and subsequently by other authors (e.g., Pejcha & Prieto 2015; Valenti et al. 2016; Müller et al. 2017), and as visible in Fig. 5, for normal SNe II there is a correlation between the ^{56}Ni mass and $M_V^{50\text{d}}$. In order to characterize the distribution of normal SNe II, we fit a straight line (solid line), where the dashed lines indicate the $\pm 2\sigma$ limits. We note that SN 1983K is below the -2σ limit, i.e., at 50 d since explosion it is brighter than those explosions producing the same amount of ^{56}Ni . The latter, along with the colour evolution of SN 1983K (Fig. 4), indicates that for that SN there is a source of photons, still relevant at 50 d since explosion, which becomes negligible at the radioactive tail. We speculate that the source of these photons is the early interaction between the SN ejecta and the CSM.

5 SPECTROSCOPIC PROPERTIES

In the previous section, we found evidence in favour of an early ejecta-CSM interaction scenario for LLEV SNe II. In this section, analysing the spectroscopic data, we present further evidence that reinforces this hypothesis.

5.1 Spectral evolution

Fig. 6a shows the early spectra of the LLEV SNe II (left), along with zooms around $\text{H}\delta$, $\text{H}\gamma$ and $\text{H}\beta$ (middle) and $\text{H}\alpha$ (right). We can see that those SNe show broad $\text{H}\alpha$ and $\text{H}\beta$ in emission, centred in the laboratory wavelengths, and with no prominent absorption features. The latter is characteristic of an SN II ejecta interacting with a CSM, where the line broadening mechanism is dominated by electron scattering (Chugai 2001). The spectral feature located between

$\text{H}\gamma$ and $\text{H}\beta$ (prominent in SN 1983K and LSQ13fn, but weaker in SN 2009aj and SN 2009au spectra) corresponds to $\text{He II } \lambda 4686$ possibly blended with $\text{N III } \lambda\lambda 4634\text{--}40\text{--}42$ (Niemela et al. 1985) and $\text{C III } \lambda 4648$ (Polshaw et al. 2016). The high temperature needed to produce the $\text{N III}/\text{C III}/\text{He II}$ feature could come from the conversion of kinetic energy into thermal energy during an ejecta-CSM interaction, which favours this scenario.

Fig. 6b shows the spectra of SN 2009au and SN 2008bm at 23 and 74 d since explosion, respectively. Since the $\text{H}\alpha$ profile in the SN 2008bm spectrum is highly contaminated by H II region lines (depicted as gray lines), we remove that contamination modelling the $\text{H}\alpha$ profile as a mixture of three Gaussians (SN $\text{H}\alpha$ absorption and emission, along with the $\text{H}\alpha$ emission from the H II region). In the middle row of the right column of Fig. 6 we see that the absorption part of the $\text{H}\alpha$ P-Cygni profile is not strong. The latter could arise from the fact that SN 2008bm and SN 2009au are fast decliners, which typically have weaker $\text{H}\alpha$ P-Cygni profile absorptions (e.g., Gutiérrez et al. 2017b), but also could indicate that the ejecta-CSM interaction is still ongoing. In Fig. 6b we also plot spectra of the fast decliners SN 2014G, SN 2013ej (e.g., Bose et al. 2015; Huang et al. 2015), and ASASSN-15oz (Bostroem et al. 2019) at epochs close in time to the SN 2009au spectrum. We see that the $\text{H}\alpha$ P-Cygni absorption component of SN 2008bm and SN 2009au is weaker (and narrower) than those of the other fast decliners, which indicates that the weakness of the $\text{H}\alpha$ P-Cygni absorption component of SN 2008bm and SN 2009au is not only related to their high photometric decline rates but also due to still ongoing ejecta-CSM interaction.

Fig. 6c shows the spectra of LSQ13fn, SN 1983K, SN 2009aj, SN 2009au, and SN 2008bm at 31, 42, 27, 27, and 80 d since explosion, respectively. At these epochs, the LLEV SNe II in our set start to show the absorption part of the $\text{H}\alpha$ P-Cygni profile, and the emission peaks appear blueshifted, where the broader the emission the more blueshifted the peak. The latter characteristic is common in normal SNe II (Anderson et al. 2014a), indicating that the effect of the ejecta-CSM interaction decreases substantially after 4–11 weeks since explosion.

Fig. 6d shows the latest spectrum of SN 2009au at 146 d since explosion, compared to spectra at similar epochs of SNe IIn (green) and SNe IIn/II (brown) that we found in the literature, and some normal SNe II (cyan). The most prominent spectral feature in the SN 2009au spectrum is the $\text{H}\alpha$ profile in emission, while the rest of the spectrum seems featureless compared to normal SNe II, as in the case of the SNe IIn and SNe IIn/II in the plot. However, we detect $[\text{Ca II}] \lambda 7291\text{--}7324$ which is present in normal SNe II but absent in the SN IIn and SN IIn/II spectra in the plot.

The left part of Fig. 7 shows the near-IR spectra of SN 2009aj combined with the closest (in time) optical spectra, along with zooms around $\text{H}\alpha$ (middle column), $\text{P}\beta$ and the spectral feature corresponding to $\text{P}\gamma$ blended with $\text{He I } \lambda 10830$ and $\text{Sr II } \lambda 10920$ (right column). For comparison, we also include combined optical and near-IR spectra of the SNe IIn (red) 2015da (Tartaglia et al. 2020) and 2010jl (Zhang et al. 2012; Borish et al. 2015), the SNe IIn/II (green) 1998S, 2007pk, and 2013fc, and the normal SN II 2017eaw (e.g., Szalai et al. 2019). All the previous SNe II were found having optical and near-IR spectra

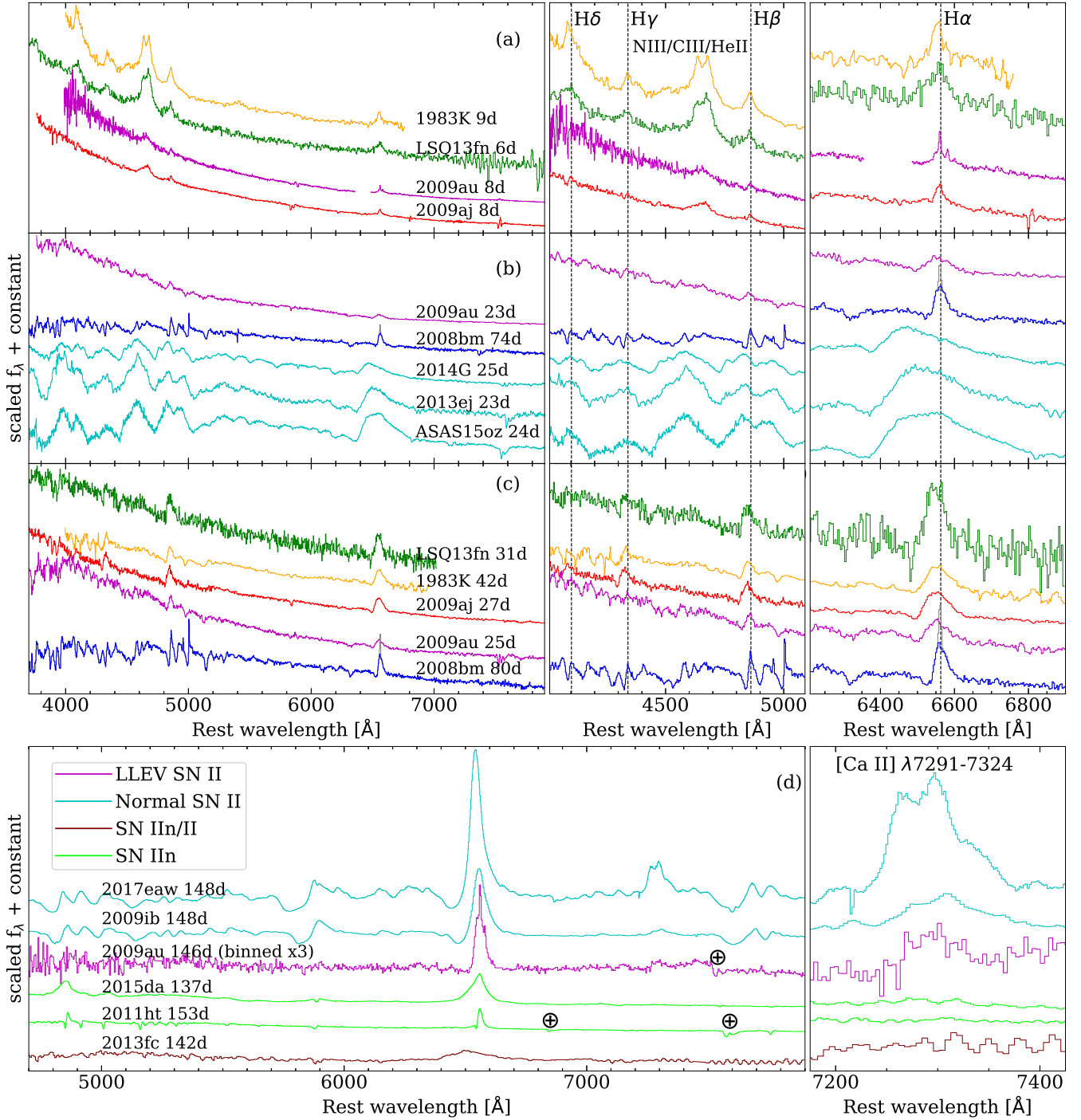


Figure 6. Panels (a), (b), and (c) show the early spectra, the last spectra still affected by the ejecta-CSM interaction, and the first spectra showing Balmer absorptions for our LLEV SN II, respectively, and zooms around $H\delta$, $H\gamma$, and $H\beta$ (middle column) and $H\alpha$ (right column). In Panel (b) we also plot normal SNe II (cyan) for comparison. Panel (d): latest spectrum of SN 2009au compared with spectra of SNe IIn, SNe IIn/II, and normal SNe II at similar epochs, and a zoom around $[\text{Ca II}] \lambda 7291\text{--}7324$ (right). Circled crosses indicate the presence of telluric lines. Epochs are since the explosion. Spectra were corrected for redshift and reddening.

close in time to the combined spectra of SN 2009aj. At 1–2 weeks since explosion, the spectrum of SN 2009aj does not show the $P\beta$ profile, being more similar to the spectrum of SN IIn 1983K than to the normal SN II 2017eaw. The absorption at $1.27 \mu\text{m}$ in the SN 2009aj spectrum is not followed by a redder emission, so this feature might not belong to the

SN. At 3–4 weeks since explosion, the spectrum of SN 2009aj is more similar to the normal SN II 2017eaw than to the SN IIn 2015da, but with $H\alpha$, $P\beta$, and $P\gamma/\text{He I}/\text{Sr II}$ profiles weaker and narrower than the observed in SN 2017eaw. At 9–10 weeks since explosion, the $H\alpha$ profile and the $P\beta$

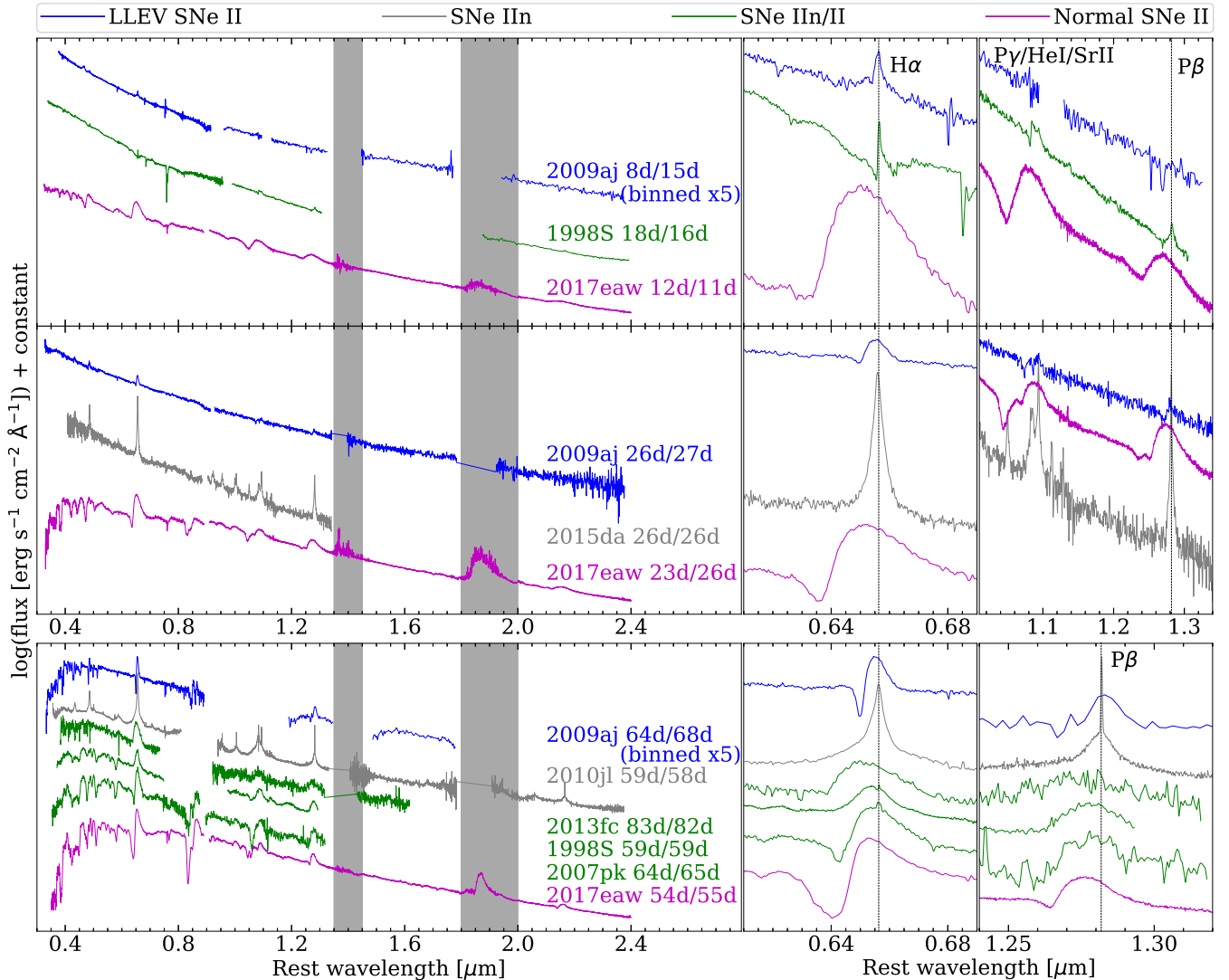


Figure 7. Left column: Near-IR spectra of SN 2009aj combined with the closest in time optical spectra (blue), compared with the combined optical and near-IR spectra of SNe IIn (gray), SNe IIn/II (green), and normal SN II 2017eaw (magenta) at a similar epochs. Middle column: zooms around $H\alpha$. Right column: zooms around the $P\gamma/HeI/SrII$ and $P\beta$ profiles.

emission component are more similar to SNe IIn/II and SN 2017eaw, but still weaker.

The aforementioned characteristics of the near-IR spectral evolution of SN 2009aj are quite similar to those found in the optical spectral evolution of the LLEV SNe II.

5.2 Weakness of metal lines

The top part of Fig. 8 shows the evolution of the $FeII \lambda 5018$ pEW of the LLEV SNe II in our set (coloured symbols), compared to the normal SNe II in the G17 sample (empty circles). We see that the LLEV SNe II are mostly below the -1σ limit (dashed line) of the G17 sample, which indicates the weakness of metal lines in the LLEV SNe II spectra in all the phase range reported in the plot.

Polshaw et al. (2016) had already reported the weakness of metal lines in the spectra of LSQ13fn compared to normal SNe II. In that work, one of the parameters that they suggested to explain the weakness of metal lines is a

low metallicity of the SN progenitor. In fact, Dessart et al. (2014) explored theoretically the dependence on the SN II progenitor metallicity of some spectral features. The general behaviour they obtained is: the lower the progenitor metallicity, the weaker the metal lines. A lower metallicity also implies a lower line blanketing and therefore a colour bluer than other SNe II with a higher metallicity, which in principle could explain the blue colour we observe in LLEV SNe II.

Nevertheless, the strength of metal lines not only depends on the progenitor metallicity, but also on the temperature of the line formation region (which modifies the opacity) and, when the SN is surrounded by a CSM, on the contribution to the flux generated from the ejecta-CSM interaction (Leloudas et al. 2015). Regarding the temperature, Anderson et al. (2016, hereafter A16) show that the pEW of the $FeII \lambda 5018$ line in normal SN II spectra increases with the $V-I$ and $V-i$ colour, which are used as a proxy for temperature.

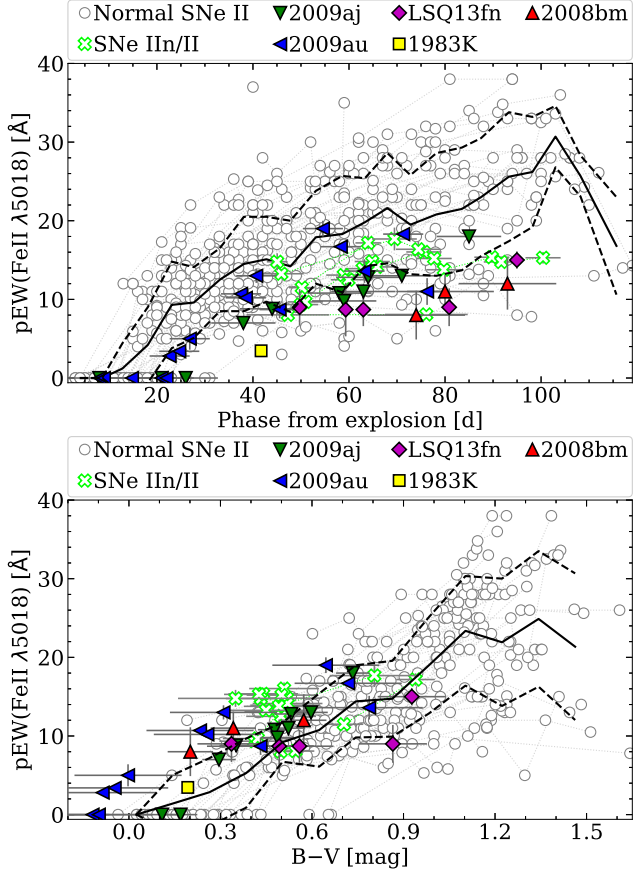


Figure 8. Evolution of the Fe II $\lambda 5018$ pEW, as a function of the time since explosion (top) and the $B-V$ colour (bottom), of the LLEV SNe II (filled symbols), the SNe II in II from the literature (empty crosses), and the normal SNe II in the G17 sample (empty circles). Black solid and dashed lines correspond to the mean values and standard deviations, respectively.

The bottom part of Fig. 8 shows the Fe II $\lambda 5018$ pEW as a function of the $B-V$ colour. In the plot we see that LLEV SNe II occupy a region similar to SNe II in II with both groups scattering around the $+1\sigma$ limit of normal SNe II. This indicates that in LLEV SNe II the weakness of metal lines is at least partially due to the higher temperature in the line forming region, but also that the line dilution play a role. Regarding the latter, to quantify its effect to first order we select SN 2008M (A14, G17) because it is one of the bluest normal SNe II in the D18 sample. With this, we can isolate as best as possible the effect of CSM dilution on the pEW of the metal lines from the similar effect of temperature we have analysed previously. Fig. 9 shows the spectrum of SN 2009aj at 39 d since explosion (red line) and a spectrum of SN 2008M nearly at the same epoch and scaled to the SN 2009aj distance (gray line). Spectra of SN 2009aj and SN 2008M have Fe II $\lambda 5018$ pEW of 7.1 and 9.7 Å, respectively. In order to include the contribution of a CSM to the flux of SN 2008M, we model the CSM as a blackbody. We choose the CSM parameters such that the continuum of SN 2008M plus the CSM matches to the continuum of SN 2009aj, which is reached using a radius of 10^{15} cm and a temperature of 7200 K. The Fe II $\lambda 5018$ pEW of the composite spectrum is of 5.7 Å, which matches better to the

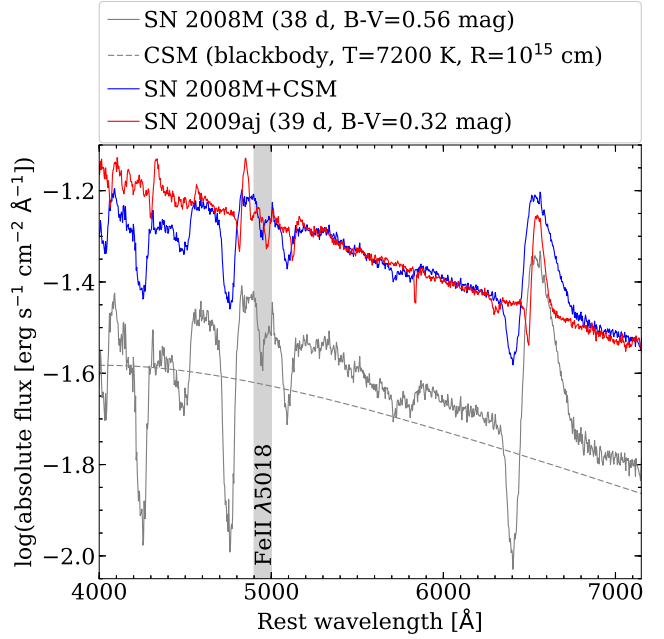


Figure 9. Comparison of the spectrum of SN 2009aj at 39 d since explosion (red line) with a composite spectrum (blue line), which is the sum of the SN 2008M spectrum at 38 d since explosion (gray line) and a blackbody continuum (dashed line). Gray region indicates the location of the Fe II $\lambda 5018$ line.

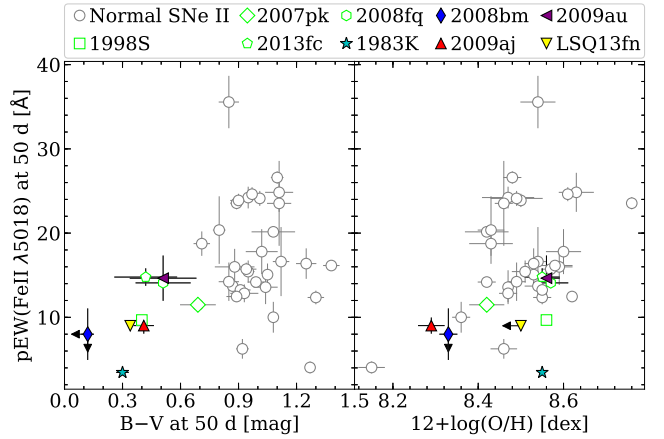


Figure 10. pEW of the Fe II $\lambda 5018$ line at 50 d since explosion as a function of the $B-V$ colour (corrected for reddening) at the same epoch (left), and as a function of the oxygen abundance (right), showing the LLEV SNe II in our set (coloured filled symbols), the SNe II in II from the literature (green empty symbols) and the normal SNe II in the A14 sample (gray empty circles). Arrows indicate upper limits.

value for SN 2009aj. Therefore, even when the SN 2008M spectrum is not as blue as the spectrum of SN 2009aj, and the oxygen abundance at the site of SN 2008M (8.43 dex; A16, which we use as a proxy for the progenitor metallicity) is higher than the value for SN 2009aj (8.29 dex; see Table 1), the flux contribution due to the ejecta-CSM interaction seems to be an efficient driver of the observed weakness of metal lines in its spectra.

Fig. 10 shows the Fe II $\lambda 5018$ pEW at 50 d since explosion, as a function of the $B-V$ colour (corrected for

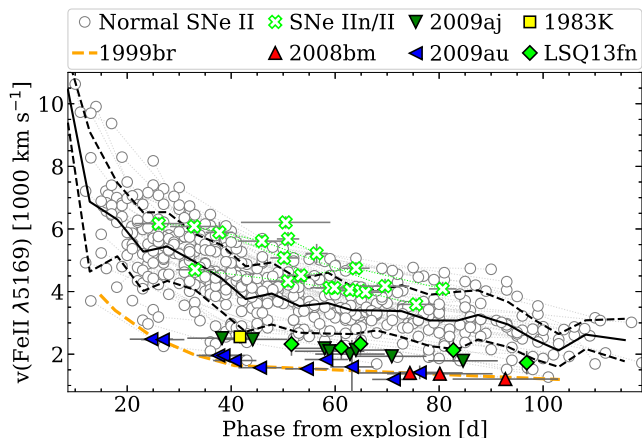


Figure 11. Evolution of the expansion velocities, measured through the Fe II $\lambda 5169$ minimum absorption lines, of the LLEV SNe II (filled symbols), the SNe IIn/II from the literature (empty crosses), and the normal SNe II in the G17 sample (empty circles). Black solid and dashed lines correspond to the mean values and standard deviations, respectively, of the expansion velocities in the G17 sample. The orange dashed line corresponds to the expansion velocities for SN 1999br.

E_{B-V}^{host}) at the same epoch (left), and as a function of the oxygen abundance in the N2 calibration of Marino et al. (2013) (right) for the LLEV SNe II in our set (coloured filled symbols). For comparison, we plot the normal SNe II in the A16 sample¹¹ (gray empty circles, where we use the colour curves plotted in Fig. 4 to compute $B-V$ colours at 50 d since explosion), and the SNe IIn/II (green empty symbols). In the left panel we see that SN 1998S, SN 2009aj, and LSQ13fn have similar Fe II $\lambda 5018$ pEW values and $B-V$ colours, while in the right panel we see that the same SNe have different oxygen abundances. The latter indicates that (1) the metallicity is not a dominant component determining the Fe II $\lambda 5018$ pEW values as $B-V$ colour is, or (2) SN 1998S and LSQ13fn suffer more line dilution compared to SN 2009aj. In Fig. 10 we also see SN 2013fc, LSQ13fn, and SN 1983K having similar $B-V$ colour and oxygen abundances, but different Fe II $\lambda 5018$ pEW values. One explanation for this difference is that SN 1983K and LSQ13fn suffer more line dilution than SN 2013fc. Based on these findings, we suggest that the progenitor metallicity is not the main driver of the weakness of metal lines seen on LLEV SNe II spectra, but a combination of higher temperatures at the line formation region and line dilution (both being consequence of an early ejecta-CSM interaction).

5.3 Expansion velocities

Fig. 11 shows the evolution of the expansion velocities, measured through the Fe II $\lambda 5169$ minimum absorption lines (v_{FeII}), of the LLEV SNe II in our set (coloured symbols) and the normal SNe II in the G17 sample¹¹ (empty circles). Expansion velocities for LLEV SNe II are all below the -1σ limit (dashed line) of the G17 sample (2820 km s^{-1} at 50 d since explosion). In particular, SN 2008bm and SN 2009au

have expansion velocities similar to SN 1999br (orange dashed line), which is one of the SNe II with the lowest expansion velocities in the G17 sample ($v_{\text{FeII}} \approx 1550 \text{ km s}^{-1}$ at 50 d since explosion). These low expansion velocities can be explained by the loss of kinetic energy by the ejecta due to their interaction with CSM.

6 DISCUSSION

6.1 A break in the magnitude-velocity relation

As we have seen in Section 4.1 and 5.3, LLEV SNe II are characterized by having luminous peak magnitudes and low expansion velocities. This is not expected in a scenario where more energetic SN II explosions that produce high luminosities also have high expansion velocities. This trend is shown in Fig. 12, where we plot the V -band absolute magnitude and the expansion velocity, both at 50 d post explosion, for the normal SNe II in the A14 SN sample (gray circles, where expansion velocities are from Table 3 of Gutiérrez et al. 2017b). To characterize the distribution of the A14 sample in this space, we perform a Gaussian process fit (solid line), where dashed lines indicate the $\pm 3\sigma$ error around the fit. We can see that the LLEV SNe II¹² (blue squares) populate a region where the only known member was LSQ13fn, marking a discrepancy of 2–3 mag with respect to the trend obtained with the A14 sample. At this point, it is not clear whether the LLEV SNe II can be considered as a separated class of SNe II, or whether they are part of a continuum in the SN II distribution induced by an increasing ejecta-CSM interaction.

To test the latter, we include in Fig. 12 the SNe IIn/II set. We see that SN 1998S, SN 2008fq, SN 2013fc, and, possibly, SN 1979C, are brighter than the 3σ limit. However, they do not fill the gap between the distribution of normal SNe II and the LLEV SNe II.

Fig. 13 shows the absolute magnitude as a function of the expansion velocity and the $B-V$ colour, all at 50 d since explosion of the LLEV SNe II in our set (blue squares). In this space we see that the LLEV SNe II seem to form a separated set of objects from the rest of normal SNe II (gray circles). With the aim to confirm this visual finding, we run the mean-shift algorithm to search for clusters in this space. We perform simulations where points are moved within their errors (assuming a normal distribution), and we find that in all of the realizations the LLEV SNe II form a separated group with respect to the rest of SNe II. Although, given the low number of LLEV SNe II, this result has to be taken with caution, the fact that these SNe always cluster in a different parameter space region, indicates that these SNe could indeed be a new sub-type of SNe II.

6.2 Interaction of the ejecta with a massive CSM

As mentioned in the previous paragraphs, a scenario where the ejecta of a normal SN II interacts with a CSM medium could explain all the characteristics observed in LLEV

¹¹ We remove SN 2008bm, SN 2009aj, and SN 2009au from the sample since they are in our LLEV SN II sample.

¹² For SN 1983K and SN 2008bm we extrapolate the expansion velocities at 50 d since explosion using the power-law given in Nugent et al. (2006).

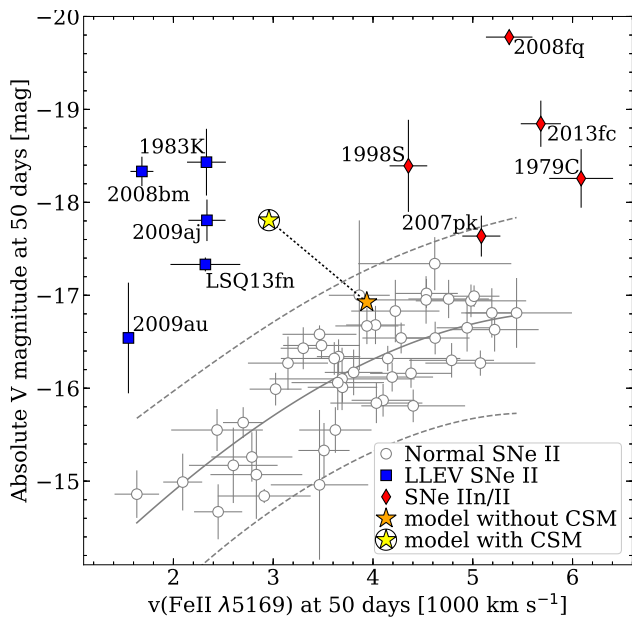


Figure 12. Absolute V -band magnitude versus expansion velocity, both at 50 d since explosion, showing the LLEV SN II set (blue squares), the normal SNe II in the A14 sample (empty circles), and the SNe IIn/II from the literature (red diamonds). The solid line corresponds to the Gaussian process fit, where dashed lines indicate the 3σ error around the fit. We also plot the location of our model for SN 2009aj (see Section 6.2) without (orange star) and with (circled yellow star) CSM.

SNe II. In fact, during the ejecta-CSM interaction, part of the kinetic energy is converted into thermal energy and photons, which slows down the ejecta, increases the temperature and luminosity, and veils the spectral lines.

In order to test whether the ejecta-CSM interaction is responsible of converting a normal SN II into a LLEV SN II, we perform hydrodynamical simulations. We focus on the modelling of SN 2009aj here because its explosion epoch and host galaxy colour excess are better constrained than those for SN 2008bm and SN 2009au. We adopt the same numerical method as in [Moriya et al. \(2017, 2018\)](#) and use the radiation hydrodynamics code STELLA ([Blinnikov et al. 1998, 2000, 2006](#)) for our numerical light curve modelling.

We take the $14 M_{\odot}$ progenitor model in [Moriya et al. \(2018\)](#) and attached a dense CSM above the progenitor by adopting the β law wind velocity. We take one model with the mass-loss rate of $3 \times 10^{-2} M_{\odot} \text{ yr}^{-1}$, the terminal wind velocity of 10 km s^{-1} , $\beta = 5$, and the dense CSM radius of 10^{15} cm . This dense CSM has a mass of $3.6 M_{\odot}$. When we explode the progenitor in this system with the explosion energy of $8 \times 10^{50} \text{ erg}$, we obtain the $UBVRI$ absolute light curves, and expansion velocities as in Fig. 14 (solid lines). We can see that the observed absolute magnitudes and expansion velocities of SN 2009aj (empty symbols) match well to the numerical results, adopting an explosion epoch 1.5 d prior to the last nondetection¹³.

The effect on the luminosity-velocity relation produced by the interaction of the ejecta with a massive CSM is de-

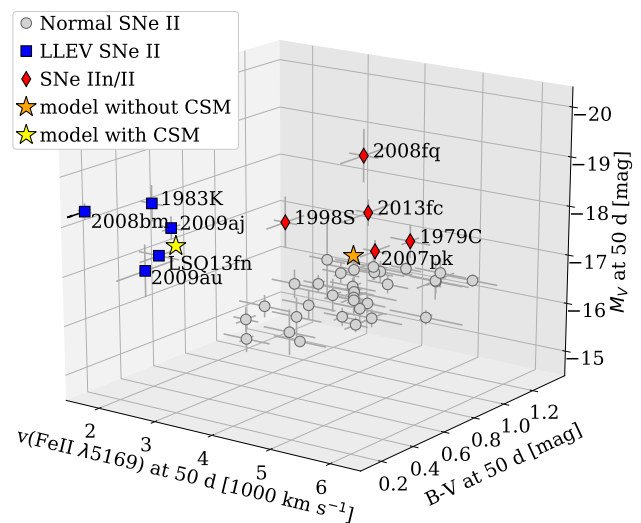


Figure 13. Absolute V -band magnitude versus expansion velocity and $B-V$ colour, all at 50 d since explosion, for the LLEV SNe II in our set (blue squares), the normal SNe II in the A14/D18 sample (gray circles), and the SNe IIn/II from the literature (red diamonds). The $B-V$ colour of SN 2008bm is an upper limit. We also plot the location of our model for SN 2009aj (see Section 6.2) without (orange star) and with (circled yellow star) CSM.

icted in Fig. 12. We can see that the SN II model without CSM (orange star), consistent with the rest of the observed SNe II, is translated outside of the 3σ limit in direction to the location of the LLEV SNe II if we include a massive CSM (circled yellow star).

6.3 The observed fraction of LLEV SNe II

In order to estimate a first approximation of the fraction of LLEV within the normal SN II family, we use the monitoring campaign carried out by the CHASE survey. Since LLEV SNe II are typically more luminous than normal SNe II, and given that the CHASE is a magnitude-limit survey, the estimation of the fraction of LLEV SNe II we will provided should be regarded as an upper limit.

During March 2008–December 2014, CHASE reported the discovery of 46 normal SNe II and 3 SNe IIn (among them, SN 2009aj and SN 2009au). In addition, we include 58 normal SNe II and 6 SNe IIn discovered by other surveys during the same period, which were discovered by CHASE independently. In order to identify possible LLEV SNe II in the CHASE normal SN II sample, we check the classification reports (and classification spectra, if they are available). We select as possible LLEV SNe II those SNe whose classification spectra best match to a subluminous SN II (e.g., SN 2005cs), or have low expansion velocities (e.g., $H_{\alpha} \lesssim 5000 \text{ km s}^{-1}$). From this list, we discard those SNe having absolute magnitude¹⁴ (at the moment of the discovery or in posterior confirmation epochs) consistent with subluminous SNe II ($\gtrsim -16.0 \text{ mag}$; [Lisakov et al. 2018](#)). From the CHASE SN IIn sample we select as possible LLEV SN II

¹³ Given the limiting magnitude of the last non-detection, we cannot discard the presence of the SN at $> 18.4 \text{ mag}$.

¹⁴ Discovery and confirmation magnitudes reported by CHASE are obtained from unfiltered images, which are similar to R -band magnitudes.

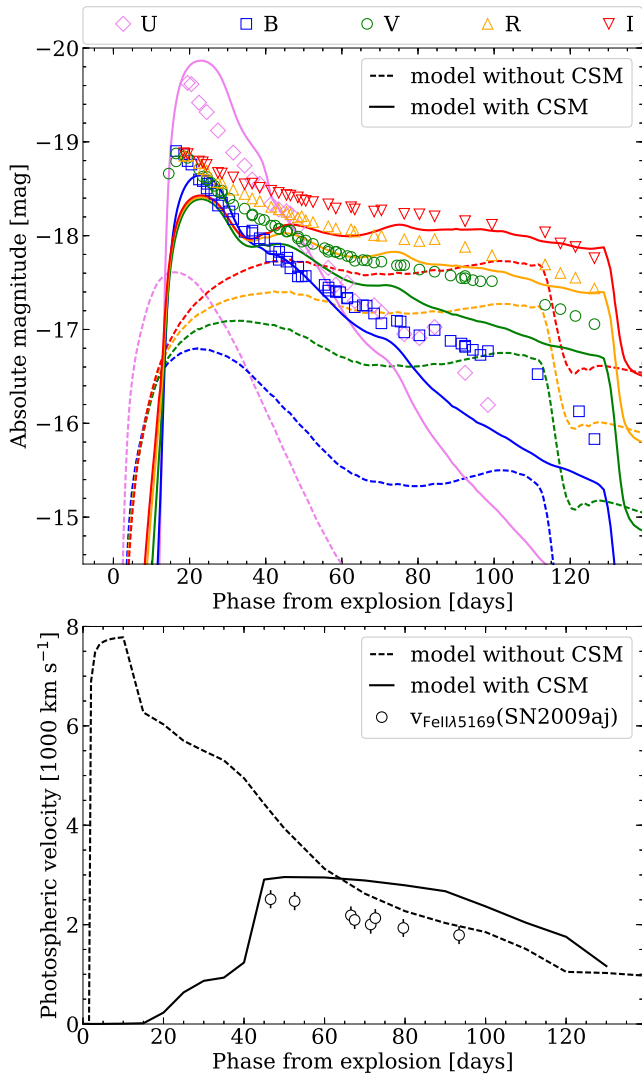


Figure 14. Absolute *UBVRI* light curves (top) and expansion velocities (middle) of our model with CSM (solid lines) and of SN 2009aj (empty symbols). The same model but without CSM (dashed lines) is displayed for comparison.

those SNe IIn that spectroscopically do not evolve as an SN IIn but as a normal SN II with narrow P-Cygni profiles. For this, we use the information available in the literature.

Among the 104 normal SNe II, 15 show spectral characteristics consistent with low expansion velocities, of which 7 are published subluminal SNe II, and 6 have absolute magnitudes (corrected for E_{B-V}^{host} inferred from Na I D at the redshift of the host galaxy) > -15.5 mag. Only SN 2009aj shows narrow P-Cygni profiles and high luminosities, while for SN 2010jc we do not have enough information to confirm or discard it. Among the 9 SNe IIn, 7 of them evolve as SNe IIn. Only SN 2009au shows a posterior evolution different than an SN IIn, while for SN 2008gm there is not enough information to confirm or discard its LLEV nature. If we consider SN 2009aj and SN 2009au as the only LLEV SNe II in the CHASE sample, then its fraction could be around 2 per cent. On the other hand, if we consider all the unconfirmed candidates (SN 2010jc and SN 2008gm) as LLEV SNe II, then the upper limit increases to 4 per cent.

The low fraction of LLEV SNe II points toward an uncommon progenitor, which should have experienced a high mass loss rate close in time to core collapse in order to generate a sufficiently massive shell ($\sim 4 M_{\odot}$) close to its surface to produce the spectroscopic and photometric characteristics that we observe in LLEV SNe II.

6.4 Impact on SNe II as distance indicators

Since the LLEV SNe II do not follow the luminosity-velocity relation observed for normal SNe II, we have to analyse their impact over the use of normal SNe II as distance indicators.

As mentioned in Rodríguez et al. (2019), the SN II distance precision using the Photospheric Magnitude Method (Rodríguez et al. 2014) could be up to 0.23 mag within a 99 per cent of confidence level. In this case, the location of the LLEV SNe II (≥ 2 mag brighter than the expected from their expansion velocities) in the Hubble diagram (HD) will be $\geq 8.7\sigma$ below the Hubble law fit. Therefore these events can be easily discarded with a sufficient amount of data.

Since the oxygen abundance of LLEV SNe II is not necessarily low, we do not have evidence that those SNe are related to low metallicity galaxies, so we do not expect that the fraction of LLEV SNe II will increase with redshift (e.g., $z > 0.5$) as consequence of the evolution of the metallicity of the Universe. However, we expect to find a higher fraction of LLEV SNe II due to the Malmquist bias. For example, for the Large Synoptic Survey Telescope (LSST) the r -band 5σ limiting magnitude is estimated to be ~ 24.3 mag, so for $z \geq 0.5$ we expected to detect SNe II with $M_r \lesssim -18.0$ mag. In this case, the observed fraction of LLEV SNe II could be as high as 10–18 per cent, so the LLEV SNe II will be only ≥ 2.9 – 3.9σ below the Hubble law, which makes it difficult to recognize them as outliers. This means that the existence of LLEV SNe II has to be taken into account if normal SNe II at high redshifts are used to derive cosmological parameters. Nevertheless, as shown in this work, the current sample have expansion velocities which are in the lower end of the normal SN II distribution, therefore a cut on expansion velocities could greatly reduce their contamination.

7 CONCLUSIONS

In this work we presented optical and near-IR data of SN 2008bm, SN 2009aj, and SN 2009au. From the analysis of these data together with already published data we found that they show similar characteristics with those of SN 1983K and LSQ13fn. In the luminosity-expansion velocity plane this possible family of SNe II, that we call LLEV, forms a separate group which have V -band absolute magnitudes 2–3 mag brighter than those expected from their expansion velocities by the luminosity-velocity relation observed for normal SNe II.

The ejecta-CSM interaction observed in LLEV SNe II lasts up to 4–11 weeks since the explosion. Subsequently, spectra show P-Cygni profiles characterized by low expansion velocities, and a weakness of metal lines. We found evidence that the metal line weakness seems not to be related to the metallicity of the progenitor but with a combined effect of the line dilution due to the contribution of the CSM to the flux, and the higher temperatures than normal SNe II at

similar epochs. Through hydrodynamic simulations, which consider a RSG progenitor of $14 M_{\odot}$ with an explosion energy of 8×10^{50} erg, we found that the high luminosity and low expansion velocities seen on SN 2009aj can be explained if the ejecta interacts with a CSM of $\sim 3.6 M_{\odot}$ located very close to the progenitor.

Based on the discoveries by the CHASE survey, we estimated an upper limit for the LLEV SNe II fraction to be 2–4 per cent of all normal SNe II. This low fraction, together with the high CSM mass we obtain from the hydrodynamic simulations, may indicate an uncommon progenitor with a high mass loss rate close in time to core collapse in order to generate a sufficiently massive shell close to its surface. Based on the available data, it is unclear whether the LLEV SNe II are a separated class of SNe II with a different progenitor system, or there is a continuum of objects connecting them with the normal SNe II. It is necessary to populate the luminosity-velocity space with more LLEV SNe II in order to reveal the nature of these peculiar SNe II. Finally, we showed that, based on the current sample, LLEV SNe II should not represent a severe contaminant in the use of normal SNe II as standardizable candles.

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APPENDIX A: TABLES AND FIGURES

This paper has been typeset from a $\text{\TeX}/\text{\LaTeX}$ file prepared by the author.

Table A1. Optical magnitudes of the sequence stars in the field of SN 2008bm. The full table is available online.

Star	α_{2000}	δ_{2000}	u (mag)	g (mag)	r (mag)	i (mag)	B (mag)	V (mag)
1	13:03:04.20	+10:28:43.0	–	–	–	14.570(11)	15.606(07)	14.907(07)
2	13:02:56.57	+10:29:47.3	16.667(25)	15.587(12)	15.120(03)	14.938(09)	15.918(08)	15.311(07)
3	13:02:50.77	+10:31:10.4	16.569(50)	15.497(50)	15.498(41)	15.567(17)	15.763(85)	15.490(56)

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

Table A2. Near-IR magnitudes of the sequence stars in the field of SN 2008bm. The full table is available online.

Star	α_{2000}	δ_{2000}	Y (mag)	J (mag)	H (mag)
1	13:03:16.69	+10:33:09.6	13.140(18)	12.883(29)	12.492(09)
2	13:02:55.18	+10:28:25.5	13.215(14)	12.788(34)	12.072(32)
3	13:03:05.57	+10:26:33.2	13.459(14)	13.071(34)	12.461(35)

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

Table A3. Optical magnitudes of the sequence stars in the field of SN 2009aj. The full table is available online.

Star	α_{2000}	δ_{2000}	u (mag)	g (mag)	r (mag)	i (mag)	z (mag)	U (mag)	B (mag)	V (mag)	R (mag)	I (mag)
1	13:56:26.43	-48:31:56.8	–	–	–	13.978(15)	–	–	14.989(32)	14.316(31)	–	–
2	13:57:02.68	-48:30:11.4	–	–	–	13.775(15)	13.739(40)	–	15.109(25)	14.259(27)	–	–
3	13:56:55.56	-48:33:32.8	–	–	–	14.081(03)	14.103(42)	–	15.127(25)	14.422(29)	–	–

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

Table A4. Near-IR magnitudes of the sequence stars in the field of SN 2009aj. The full table is available online.

Star	α_{2000}	δ_{2000}	Y (mag)	J (mag)	H (mag)
1	13:56:56.40	-48:29:39.5	12.653(10)	12.400(10)	12.092(11)
2	13:56:47.70	-48:33:36.7	12.501(06)	12.189(07)	11.754(07)
3	13:56:43.43	-48:25:06.7	13.101(07)	12.896(12)	12.648(18)

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

Table A5. Optical magnitudes of the sequence stars in the field of SN 2009au. The full table is available online.

Star	α_{2000}	δ_{2000}	u (mag)	g (mag)	r (mag)	i (mag)	z (mag)	B (mag)	V (mag)	R (mag)	I (mag)
1	12:59:31.60	-29:32:03.1	16.470(35)	15.097(22)	14.621(11)	14.400(21)	14.327(20)	15.513(30)	14.801(26)	14.414(22)	13.983(22)
2	12:59:51.12	-29:37:51.5	17.176(107)	15.316(26)	14.694(11)	14.437(22)	14.329(16)	15.805(25)	14.931(22)	14.467(22)	14.000(26)
3	12:59:38.00	-29:32:02.3	16.256(48)	14.940(24)	14.483(11)	14.236(09)	14.155(18)	15.349(27)	14.654(25)	14.279(24)	13.822(24)

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

Table A6. Near-IR magnitudes of the sequence stars in the field of SN 2009au. The full table is available online.

Star	α_{2000}	δ_{2000}	Y (mag)	J (mag)	H (mag)
1	13:00:04.82	-29:38:23.8	12.453(21)	11.976(15)	11.317(09)
2	13:00:10.53	-29:35:09.2	12.587(22)	12.206(18)	11.675(12)
3	12:59:47.94	-29:39:32.1	12.863(17)	12.596(14)	12.309(09)

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

Table A7. CSP-I photometry of SN 2008bm. The full table is available online.

MJD	u (mag)	g (mag)	r (mag)	i (mag)	B (mag)	V (mag)	Y (mag)	J (mag)	H (mag)
54558.23	18.548(24)	17.701(09)	17.536(08)	17.579(12)	17.945(10)	17.763(10)	–	–	–
54559.34	–	–	–	–	–	–	17.257(15)	17.125(21)	17.096(59)
54560.25	18.755(32)	17.757(08)	17.584(08)	17.602(09)	18.023(11)	17.764(09)	–	–	–

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

Table A8. Photometry of SN 2009aj. The full table is available online.

MJD	u (mag)	g (mag)	r (mag)	i (mag)	z (mag)	U (mag)	B (mag)	V (mag)	R (mag)	I (mag)	Y (mag)	J (mag)	H (mag)	Telescope
54873.20	–	–	–	–	–	–	–	>18.412 ^a	–	–	–	–	–	PROMPT
54886.26	–	–	–	–	–	–	–	15.631(29) ^a	–	–	–	–	–	PROMPT
54888.18	–	–	–	–	–	–	–	15.417(31) ^a	–	–	–	–	–	PROMPT

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

^aUnfiltered photometry calibrated to the V -band.

Table A9. Photometry of SN 2009au. The full table is available online.

MJD	<i>u</i>	<i>g</i>	<i>r</i>	<i>i</i>	<i>z</i>	<i>B</i>	<i>V</i>	<i>R</i>	<i>I</i>	<i>Y</i>	<i>J</i>	<i>H</i>	Telescope
	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	(mag)	
54886.18	–	–	–	–	–	–	>19.042 ^a	–	–	–	–	–	PROMPT
54893.15	–	–	–	–	–	–	>17.656 ^a	–	–	–	–	–	PROMPT
54901.17	–	–	–	–	–	–	17.629(77) ^a	–	–	–	–	–	PROMPT

Note: Errors, in parenthesis and in units of 0.001 mag, are 1σ .

^aUnfiltered photometry calibrated to the *V*-band.

Table A10. Optical spectroscopy of SN 2009au.

UT date	MJD	Phase ^a	Instrument setup ^b	Wavelength range	Exposure time	Slit
		(d)		(Å)	(s)	(")
2009-03-22	54912.11	14.8	TNG+LRS (+LR-B, LR-R)	3320–10360	1200, 1200	1.5
2009-03-30	54920.03	22.6	TNG+LRS (+LR-B, LR-R)	3800–9190	1800, 1800	1.5
2009-04-01	54922.05	24.6	NOT+ALFOSC (+gm4)	3420–9110	900	1.0
2009-04-14	54935.04	37.5	NOT+ALFOSC (+gm4)	3490–8000	1200	1.0
2009-04-15	54936.05	38.5	NOT+ALFOSC (+gm4)	3640–9070	2700	1.3
2009-05-05	54956.01	58.3	NOT+ALFOSC (+gm4)	3700–8000	1800	1.0
2009-05-09	54960.99	63.2	NOT+ALFOSC (+gm4)	3640–9080	900	1.3
2009-05-18	54969.13	71.2	NTT+EFOSC (+gr11)	3340–10290	1200	1.0
2009-07-16	55028.08	129.6	NTT+EFOSC (+gr16)	6000–9400	2400	1.5
2009-08-01	55044.98	146.4	Gemini-S+GMOS (+R400)	4700–8120	1800	1.0

^aSince explosion (MJD 54897.2).

^bTNG+LRS: 3.58 m Telescopio Nazionale Galileo (La Palma, Spain) + Low Resolution Spectrograph; NOT+ALFOSC: 2.56 m Nordic Optical Telescope (La Palma, Spain) + Alhambra Faint Object Spectrograph and Camera; NTT+EFOSC: 3.58 m New Technology Telescope (La Silla, Chile) + ESO Faint Object Spectrograph and Camera; Gemini-S+GMOS: 8.1 m Gemini South Telescope (Cerro Pachón, Chile) + Gemini Multi-Object Spectrograph.

Table A11. Near-IR spectroscopy of SN 2009aj.

UT date	MJD	Phase ^a	Instrument setup ^b	Wavelength range	Exposure time	Slit
		(d)		(Å)	(s)	(")
2009-03-04	54895.38	15.4	VLT+ISAAC (+SWS1-LR)	9700–25000	600	1.0
2009-03-16	54907.36	27.3	NTT+SOFI (+GB, GR)	9400–24400	2700, 5400	0.6
2009-04-26	54948.17	67.7	SOAR+OSIRIS (+LR)	11700–19400	2400, 2400	1.0

^aSince explosion (MJD 54879.8).

^bVLT+ISAAC: 8.2 m Very Large Telescope UT1 (Paranal, Chile) + Infrared Spectrometer And Array Camera; NTT+SOFI: 3.58 m New Technology Telescope (La Silla, Chile) + Son of ISAAC; SOAR+OSIRIS: 4.1 m Southern Astrophysical Research Telescope (Cerro Pachón, Chile) + Ohio State Infrared Imager/Spectrometer.

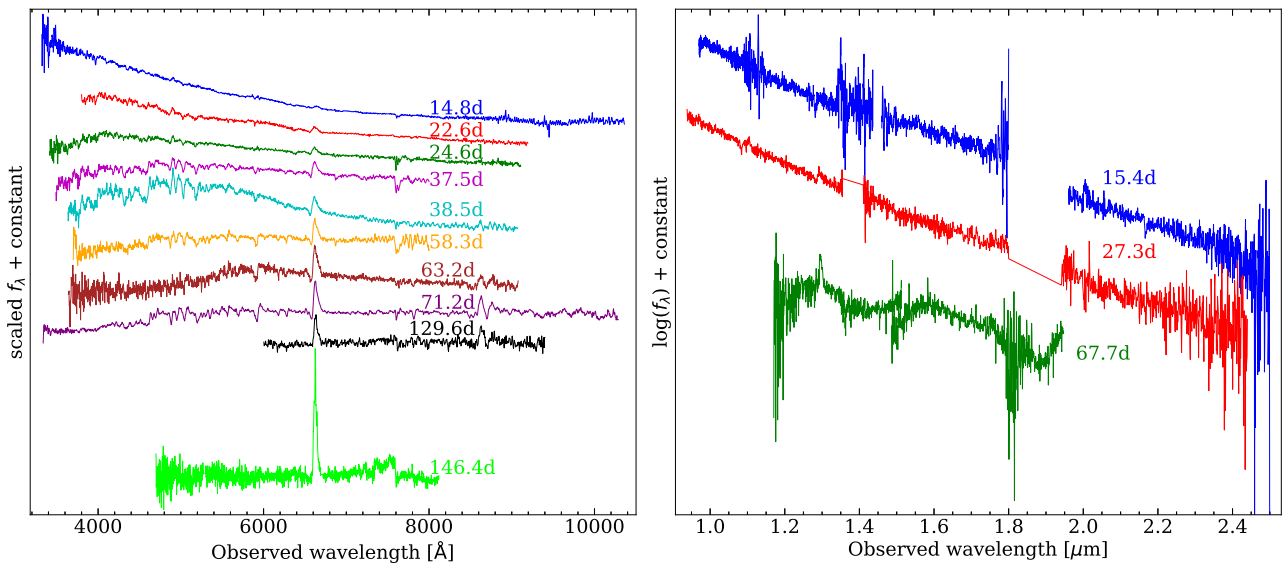
**Figure A1.** Left: optical spectra of SN 2009au. Right: near-IR spectra of SN 2009aj.

Table A12. SN and host galaxy parameters of the SN IIn/II sample.

SN data	1979C	1998S	2007pk	2008fq	PTF11iqb	2013fc
Host galaxy	M100	NGC 3877	NGC 579	NGC 6907	NGC 151	ESO 154-G10
Host type ^a	SAB(s)bc	SA(s)c?	Scd?	SB(s)bc	SB(r)bc	(R')SB(r)a?
E_{B-V}^{MW} (mag) ^b	0.023	0.020	0.045	0.055	0.028	0.026
E_{B-V}^{host} (mag)	$0.13^{+0.06}_{-0.02}$	$0.20^{+0.11}_{-0.08}$	0.09 ± 0.04^c	0.99 ± 0.14^d	0.0	0.91 ± 0.06
Explosion epoch (MJD)	43970.4 ± 8.4	50871.2 ± 3.5	54411.8 ± 2.5	54719.8 ± 4.5	55762.2 ± 2.7	56516.2 ± 2.0
cz_{helio} (km s ⁻¹) ^{a,e}	1571	895	4993	3182	3747	5586
μ_z (mag) ^f	30.74 ± 0.81	31.24 ± 0.64	34.14 ± 0.17	33.37 ± 0.24	33.51 ± 0.23	34.43 ± 0.15
μ_{zi} (mag) ^g	$31.18 \pm 0.05^{\ddagger}$	$30.92 \pm 0.45^{\dagger}$	–	$32.80 \pm 0.40^{\dagger}$	$33.89 \pm 0.40^{\dagger}$	–
Adopted μ (mag) ^h	31.18 ± 0.05	31.03 ± 0.37	34.14 ± 0.17	33.22 ± 0.21	33.60 ± 0.21	34.43 ± 0.15
$12 + \log(\text{O}/\text{H})$ (dex) ⁱ	–	$8.56 \pm 0.10^*$	$8.42 \pm 0.04^\circ$	$8.57 \pm 0.04^\otimes$	$8.61 \pm 0.03^*$	$8.55 \pm 0.04^\oplus$
s_2 (mag (100 d) ⁻¹)	3.05 ± 0.08	3.62 ± 0.09	2.28 ± 0.05	1.77 ± 0.08	1.40 ± 0.10	3.59 ± 0.12
M_V^{max} (mag)	-19.55 ± 0.14	-19.46 ± 0.49	-18.54 ± 0.21	-20.82 ± 0.50	-18.34 ± 0.23	-20.13 ± 0.24
M_V^{50d} (mag)	-18.26 ± 0.31	-18.39 ± 0.49	-17.64 ± 0.21	-19.78 ± 0.50	-17.39 ± 0.25	-18.85 ± 0.24
⁵⁶ Ni mass (M _⊙)	$> 0.069 \pm 0.010^j$	$> 0.155 \pm 0.069^j$	–	–	0.029 ± 0.007	$> 0.279 \pm 0.070^j$
$(B-V)_{50d}$ (mag)	0.53 ± 0.05	0.40 ± 0.11	0.69 ± 0.09	0.51 ± 0.14	0.46 ± 0.07	0.42 ± 0.16
pEW _{FeIIλ5018} ^{50d} (Å)	8.03 ± 0.40	9.68 ± 0.31	11.49 ± 0.14	14.09 ± 0.46	–	14.78 ± 0.98
v_{FeII}^{50d} (km s ⁻¹)	6086 ± 315	4355 ± 183	5086 ± 185	5365 ± 228	–	5682 ± 196
References ^k	1, 2, 3, 4, 5	6, 7, 8	9, 10	11, 12	13	14

^aFrom NED.

^bGalactic colour excesses from [Schlafly & Finkbeiner \(2011\)](#), with a statistical uncertainty of 16% ([Schlegel et al. 1998](#)).

^cAverage of the E_{B-V}^{host} values reported in [Pritchard et al. \(2012\)](#) and [Inserra et al. \(2013\)](#).

^dSN 2008fq has a NaID pEW of 2.9 ± 0.2 Å ([Taddia et al. 2013](#)), which is unhelpful to estimate E_{B-V}^{host} (see Section 2.3). As for SN 2009au, we compute E_{B-V}^{host} matching the $B-V$ colour curve of SN 2008fq to the rest of SNe IIn/II.

^eHeliocentric velocities, with an error of 162 km s⁻¹ to take into account the host galaxy rotational velocity ([Anderson et al. 2014a](#)).

^fDistance moduli computed from recession velocities, corrected for the infall of the Local Group towards the Virgo cluster and the Great Attractor ([Mould et al. 2000](#)), and assuming $H_0 = 73$ km s⁻¹ Mpc⁻¹. We include an error of 382 km s⁻¹ due to peculiar velocities ([Wang et al. 2006](#)).

^gRedshift-independent distance moduli (when available): (‡) Cepheids distance from [Saha et al. \(2006\)](#); (†) Tully-Fisher distance from the Extragalactic Distance Database (EDD, <http://edd.ifa.hawaii.edu/>).

^hWeighted average of μ_z and μ_{zi} (if it is available). For SN 1979C we adopt the Cepheid distance modulus.

ⁱOxygen abundances, in the N2 calibration of [Marino et al. \(2013\)](#), measured by [Anderson et al. \(2016\)](#) (⊗), [Taddia et al. \(2015\)](#) (*), [Inserra et al. \(2013\)](#) (◊), and [Kangas et al. \(2016\)](#) (⊕). The latter three are recalibrations of the original values reported in the N2 calibration of [Pettini & Pagel \(2004\)](#).

^j⁵⁶Ni masses measured with the [Hamuy \(2003\)](#) relation and $BC = 0.26 \pm 0.06$ ([Hamuy 2001](#)). The upper limits are because the slope during the radioactive tail is > 1.5 mag (100 d)⁻¹.

^kSince $s_3 > 1.5$ mag (100 d)⁻¹, ⁵⁶Ni mass estimations are lower limits.

^k(1) [Balinskaia et al. \(1980\)](#); (2) [de Vaucouleurs et al. \(1981\)](#); (3) [Branch et al. \(1981\)](#); (4) [Barbon et al. \(1982\)](#); (5) [Penston & Blades \(1980\)](#); (6) [Fassia et al. \(2000\)](#); (7) [Poon et al. \(2011\)](#); (8) [Fassia et al. \(2001\)](#); (9) [Inserra et al. \(2013\)](#); (10) [Hicken et al. \(2017\)](#); (11) [Taddia et al. \(2013\)](#); (12) [Faran et al. \(2014\)](#); (13) [Smith et al. \(2015\)](#); (14) [Kangas et al. \(2016\)](#).

Table A13. Properties of the normal SNe II with $B-V$ colours during the radioactive tail.

SN	t_0 (MJD)	$E_{B-V}^{MW \dagger}$ (mag)	E_{B-V}^{host} (mag)	$cz_{\text{helio}}^{\ddagger}$ (km s ⁻¹)	Ref* [*]	SN	t_0 (MJD)	$E_{B-V}^{MW \dagger}$ (mag)	E_{B-V}^{host} (mag)	$cz_{\text{helio}}^{\ddagger}$ (km s ⁻¹)	Ref* [*]
1969L	40550.0 ± 5.0	0.053	0.000 ± 0.097	518	1, 2	2007it	54348.0 ± 1.0	0.099	0.019 ± 0.013	1193	15, 16
1992H	48660.5 ± 10.0	0.015	0.000 ± 0.097	1793	2, 3	2008bk	54540.9 ± 8.0	0.017	0.000 ± 0.016	230	7, 16, 17
1996W	50179.5 ± 3.0	0.036	0.187	1617	4	2008gz	54693.5 ± 5.0	0.036	0.030 ± 0.040	1862	18
1999ca	51277.5 ± 7.0	0.094	0.039 ± 0.100	2791	5, 6, 7	2009bw	54916.0 ± 3.0	0.197	0.080	1155	19
1999em	51475.2 ± 3.8	0.035	0.100 ± 0.052	800	5, 6, 8	2012A	55928.7 ± 4.7	0.027	0.009	753	20
2003B	52621.7 ± 4.3	0.023	0.000 ± 0.081	1141	5, 6	2012aw	56002.1 ± 0.8	0.024	0.046 ± 0.008	778	21, 22
2003gd	52715.0 ± 3.0	0.060	0.130 ± 0.100	657	5, 6, 9	2014cx	56901.9 ± 0.5	0.096	0.0	1646	23
2004dj	53186.5 ± 3.0	0.034	0.161 ± 0.081	221	6, 10, 11	ASAS14jb	56945.6 ± 3.0	0.015	0.0	1808	24
2004et	53270.5 ± 0.5	0.293	0.000 ± 0.081	40	6, 12	ASAS15oz	57261.1 ± 4.0	0.078	0.0	2078	25
2005cs	53548.4 ± 0.5	0.032	0.040 ± 0.050	463	13, 14	2017eaw	57886.2 ± 1.0	0.293	0.0	40	26, 27, 28

[†]Galactic colour excesses from [Schlafly & Finkbeiner \(2011\)](#), with a statistical error of 16 per cent ([Schlegel et al. 1998](#)).

[‡]Heliocentric redshifts from NED.

^{*}References: (1) [Ciatti et al. \(1971\)](#); (2) [Hamuy \(2003\)](#); (3) [Clocchiatti et al. \(1996\)](#); (4) [Inserra et al. \(2013\)](#); (5) [Galbany et al. \(2016\)](#); (6) [Olivares E. et al. \(2010\)](#); (7) [Gutiérrez et al. \(2017a\)](#); (8) [Elmhamdi et al. \(2003\)](#); (9) [Van Dyk et al. \(2003\)](#); (10) [Vinkó et al. \(2006\)](#); (11) [Tsvetkov et al. \(2008\)](#); (12) [Maguire et al. \(2010\)](#); (13) [Pastorello et al. \(2009\)](#); (14) [Dessart et al. \(2008\)](#); (15) [Andrews et al. \(2011\)](#); (16) [Anderson et al. \(2014b\)](#); (17) [Van Dyk et al. \(2012\)](#); (18) [Roy et al. \(2011\)](#); (19) [Inserra et al. \(2012\)](#); (20) [Tomasella et al. \(2013\)](#); (21) [Bose et al. \(2013\)](#); (22) [Dall’Ora et al. \(2014\)](#); (23) [Huang et al. \(2016\)](#); (24) [Meza et al. \(2019\)](#); (25) [Bostroem et al. \(2019\)](#); (26) [Tsvetkov et al. \(2018\)](#); (27) [Szalai et al. \(2019\)](#); (28) [Van Dyk et al. \(2019\)](#).

Table A14. Properties and references of the SNe II used in this work.

SN	t_0 (MJD)	$E_{B-V}^{\text{MW}\dagger}$ (mag)	E_{B-V}^{host} (mag)	$cz_{\text{helio}}^{\ddagger}$ (km s $^{-1}$)	References*
1994Y	49570.1 \pm 5.9	0.008	0.0	2558	1, 2
1996al	50255.0 \pm 2.0	0.010	0.10 \pm 0.05	1970	3
2010jl	55473.5 \pm 5.0	0.023	0.02	3214	4, 5
2011ht	55833.2	0.072	0.0	1093	6
2015bh	57154.8 \pm 3.9	0.019	0.21 \pm 0.07	1947	7, 8
2015da	57030.4 \pm 1.5	0.012	0.97 \pm 0.27	2165	9

\dagger Galactic colour excesses from [Schlafly & Finkbeiner \(2011\)](#), with a statistical error of 16 per cent ([Schlegel et al. 1998](#)).

\ddagger Heliocentric redshifts from NED.

* (1) [Ho et al. \(2001\)](#); (2) [Tsvetkov & Pavlyuk \(1997\)](#); (3) [Benetti et al. \(2016\)](#); (4) [Fransson et al. \(2014\)](#); (5) [Chandra et al. \(2015\)](#); (6) [Humphreys et al. \(2012\)](#); (7) [Elias-Rosa et al. \(2016\)](#); (8) [Thöne et al. \(2017\)](#); (9) [Tartaglia et al. \(2020\)](#).

Table A15. New estimation of the nickel masses for the SNe II in the Hamuy sample.

SN*	$E_{B-V}^{\text{MW}\dagger}$ (mag)	$E_{B-V}^{\text{host}*}$ (mag)	μ_z^\diamond (mag)	μ_{zi}^\diamond (mag)	μ^\ddagger (mag)	$M_V^{50\text{d}}$ (mag)	$M(^{56}\text{Ni})$ (M_\odot)
1969L	0.053	0.000 \pm 0.097 ^a	29.78 \pm 1.26	29.84 \pm 0.04 ^c	29.84 \pm 0.04	-16.65 \pm 0.31	0.068 \pm 0.021
1970G	0.008	0.000 \pm 0.097 ^a	28.96 \pm 1.83	29.14 \pm 0.05 ^d	29.14 \pm 0.05	-17.06 \pm 0.34	0.030 \pm 0.010
1973R	0.028	0.452 \pm 0.097 ^a	29.07 \pm 1.74	30.18 \pm 0.04 ^e	30.18 \pm 0.04	-17.11 \pm 0.31	0.093 \pm 0.036
1986I	0.034	0.065 \pm 0.097 ^a	33.00 \pm 0.29	30.71 \pm 0.40 ^f	30.71 \pm 0.40	-16.47 \pm 0.54	0.076 \pm 0.036
1988A	0.035	0.000 \pm 0.097 ^a	30.71 \pm 0.82	31.11 \pm 0.40 ^f	31.03 \pm 0.36	-16.14 \pm 0.47	0.054 \pm 0.027
1990E	0.022	0.468 \pm 0.097 ^a	31.07 \pm 0.69	30.89 \pm 0.45 ^f	30.94 \pm 0.38	-16.56 \pm 0.52	0.044 \pm 0.022
1990K	0.012	0.065 \pm 0.097 ^a	31.66 \pm 0.53	31.82 \pm 0.45 ^f	31.75 \pm 0.34	-17.49 \pm 0.50	0.036 \pm 0.017
1991G	0.017	0.000 \pm 0.097 ^a	30.90 \pm 0.75	30.64 \pm 0.45 ^f	30.71 \pm 0.39	-15.23 \pm 0.50	0.020 \pm 0.009
1992H	0.015	0.000 \pm 0.097 ^a	32.53 \pm 0.35	32.23 \pm 0.45 ^f	32.42 \pm 0.28	-17.48 \pm 0.41	0.139 \pm 0.054
1992ba	0.050	0.139 \pm 0.052 ^b	30.76 \pm 0.80	–	30.76 \pm 0.80	-15.92 \pm 0.82	0.024 \pm 0.018
1999gi	0.014	0.181 \pm 0.052 ^b	30.01 \pm 1.13	30.77 \pm 0.06 ^c	30.77 \pm 0.06	-16.47 \pm 0.18	0.041 \pm 0.007

Note. To recompute $M_V^{50\text{d}}$ and $M(^{56}\text{Ni})$ we use the same observables listed in Table 2 of [Hamuy \(2003\)](#).

*We recompute $M_V^{50\text{d}}$ and $M(^{56}\text{Ni})$ values only for the SNe that do not have $M(^{56}\text{Ni})$ estimations in the A14 sample, with three or more V -band photometric points during the nebular phase, and with $s_3 < 1.3$ mag (100 d) $^{-1}$.

\dagger Galactic colour excesses from [Schlafly & Finkbeiner \(2011\)](#), with a statistical error of 16 per cent ([Schlegel et al. 1998](#)).

*Host galaxy reddening values from [Hamuy \(2003\)](#) (a) or [Olivares E. et al. \(2010\)](#) (b).

\diamond Distance moduli computed from recession velocities, corrected for the infall of the Local Group towards the Virgo cluster and the Great Attractor ([Mould et al. 2000](#)), and assuming $H_0 = 73$ km s $^{-1}$ Mpc $^{-1}$. We include an error of 382 km s $^{-1}$ due to peculiar velocities ([Wang et al. 2006](#)).

\circ Redshift-independent distance moduli: Cepheid distance from [Saha et al. \(2006\)](#) (c) and [Riess et al. \(2016\)](#) (d); TRGB distance from [Jang & Lee \(2017\)](#) (e); Tully-Fisher distance from EDD (f). We do not consider distances computed with the SN itself.

\ddagger Weighted average of μ_z and μ_{zi} (if it is available). For SN 1969L, SN 1970G, SN 1973R, SN 1986I, and SN 1999gi we adopt the redshift-independent distance modulus.