

# q-Fourier Transform and its inversion-problem

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## Abstract

Tsallis' q-Fourier transform is not generally one-to-one. It is shown here that, if we eliminate the requirement that  $q$  be fixed, and let it instead "float", a simple extension of the  $F_q$ -definition, this procedure

restores the one-to-one character.

KEYWORDS:  $q$ -Fourier transform, generalization, one to one character, statistical mechanics, nonextensive statistical mechanics.

# 1 Introduction

Nonextensive statistical mechanics (NEXT) [1, 2, 3], a current generalization of the BoltzmannGibbs (BG) one, is actively studied in diverse areas of Science. NEXT is based on a nonadditive (though extensive [4]) entropic information measure characterized by the real index  $q$  (with  $q = 1$  recovering the standard BG entropy). It has been applied to variegated systems such as cold atoms in dissipative optical lattices [5], dusty plasmas [6], trapped ions [7], spinglasses [8], turbulence in the heliosheath [9], self-organized criticality [10], high-energy experiments at LHC/CMS/CERN [11] and RHIC/PHENIX/Brookhaven [12], low-dimensional dissipative maps [13], finance [14], galaxies [15], Fokker-Planck equation's applications [16], etc.

NEXT can be advantageously expressed via  $q$ -generalizations of standard mathematical concepts (the logarithm and exponential functions, addition and multiplication, Fourier transform (FT) and the Central Limit Theorem (CLT) [17, 22, 25]). The  $q$ -Fourier transform  $F_q$  exhibits the nice property of transforming  $q$ -Gaussians into  $q$ -Gaussians [17]. Recently, plane waves, and the representation of the Dirac delta into plane waves have been also generalized [18, 19, 21, 22].

A serious problem afflicts  $F_q$ . It is not generally one-to-one. A detailed

example is discussed below. In this work we show that by recourse to a rather simple but efficient stratagem that consists in

- eliminating the requirement that  $q$  be fixed and instead
- let it “float”,

one restores the one-to-one character.

## 2 Generalizing the $q$ -Fourier transform

We define, following [17], a  $q$ -Fourier transform of  $f(x) \in L^1(\mathbb{R})$ ,  $f(x) \geq 0$  as

$$F(k, q) = [H(q - 1) - H(q - 2)] \times \int_{-\infty}^{\infty} f(x) \{1 + i(1 - q)kx[f(x)]^{(q-1)}\}^{\frac{1}{1-q}} dx \quad (2.1)$$

where  $H(x)$  is the Heaviside step function.

The only difference between this definition and that given in [17] is that  $q$  is not fixed and varies within the interval  $[1, 2)$ . Herein lies the hard-core of our presentation. This simple change of perspective makes it is easy to find the inversion-formula for (2.1) by recourse to the inverse Fourier transform

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \left[ \lim_{\epsilon \rightarrow 0^+} \int_1^2 F(k, q) \delta(q - 1 - \epsilon) dq \right] e^{-ikx} dk. \quad (2.2)$$

As a consequence, we see that this  $q$ -Fourier transform is one-to-one, unlike what happens in [23],[24]. In the next section we give an illustrative example.

### 3 Example

As an illustration we discuss the example given by Hilhorst in Ref. ([22]).

Let  $f(x)$  be

$$f(x) = \begin{cases} \left(\frac{\lambda}{x}\right)^\beta ; & x \in [a, b] ; 0 < a < b ; \lambda > 0 \\ 0 ; & x \text{ outside } [a, b]. \end{cases} \quad (3.1)$$

The corresponding  $q$ -Fourier transform is

$$F(k, q) = \lambda^\beta \int_a^b x^{-\beta} \{1 + i(1 - q)k\lambda^{\beta(q-1)}x^{1-\beta(q-1)}\}^{\frac{1}{1-q}} dx. \quad (3.2)$$

Effecting the change of variables

$$y = x^{1-\beta(q-1)},$$

we have for (3.2)

$$F(k, q) = [H(q - 1) - H(q - 2)] \times \frac{\lambda^\beta}{1 - \beta(q - 1)} \int_{a^{1-\beta(q-1)}}^{b^{1-\beta(q-1)}} y^{\frac{\beta(q-2)}{1-\beta(q-1)}} \{1 + i(1 - q)k\lambda^{\beta(q-1)}y\}^{\frac{1}{1-q}} dy. \quad (3.3)$$

Now, (3.3) can be rewritten in the useful form

$$\begin{aligned}
F(k, q) &= [H(q-1) - H(q-2)] \times \\
&\left\{ \left\{ H(q-1) - H \left[ q - \left( 1 + \frac{1}{\beta} \right) \right] \right\} \times \right. \\
&\frac{\lambda^\beta}{1 - \beta(q-1)} \int_{a^{1-\beta(q-1)}}^{b^{1-\beta(q-1)}} y^{-\frac{\beta(2-q)}{1-\beta(q-1)}} \{1 + i(1-q)k\lambda^{\beta(q-1)}y\}^{\frac{1}{1-q}} dy + \\
&\left. \left\{ H \left[ q - \left( 1 + \frac{1}{\beta} \right) \right] - H(q-2) \right\} \times \right. \\
&\left. \frac{\lambda^\beta}{\beta(q-1) - 1} \int_{b^{1-\beta(q-1)}}^{a^{1-\beta(q-1)}} y^{\frac{\beta(q-2)}{1-\beta(q-1)}} \{1 + i(1-q)k\lambda^{\beta(q-1)}y\}^{\frac{1}{1-q}} dy \right\}. \quad (3.4)
\end{aligned}$$

Taking into account that the involved integrals are defined in a finite interval,

we can cast (3.4) as

$$\begin{aligned}
F(k, q) &= [H(q-1) - H(q-2)] \times \\
&\left\{ \left\{ H(q-1) - H \left[ q - \left( 1 + \frac{1}{\beta} \right) \right] \right\} \times \right. \\
&\frac{\lambda^\beta}{1 - \beta(q-1)} \lim_{\epsilon \rightarrow 0^+} \int_{a^{1-\beta(q-1)}}^{b^{1-\beta(q-1)}} y^{-\frac{\beta(2-q)}{1-\beta(q-1)}} \{1 + i(1-q)(k+i\epsilon)\lambda^{\beta(q-1)}y\}^{\frac{1}{1-q}} dy + \\
&\left. \left\{ H \left[ q - \left( 1 + \frac{1}{\beta} \right) \right] - H(q-2) \right\} \times \right. \\
&\left. \frac{\lambda^\beta}{\beta(q-1) - 1} \lim_{\epsilon \rightarrow 0^+} \int_{b^{1-\beta(q-1)}}^{a^{1-\beta(q-1)}} y^{\frac{\beta(q-2)}{1-\beta(q-1)}} \{1 + i(1-q)(k+i\epsilon)\lambda^{\beta(q-1)}y\}^{\frac{1}{1-q}} dy \right\}. \quad (3.5)
\end{aligned}$$

We now use results of the Integral's table [26] to evaluate (3.5) and get

$$\begin{aligned}
& \lim_{\epsilon \rightarrow 0^+} \int_{a^{1-\beta(q-1)}}^{\infty} y^{-\frac{\beta(2-q)}{1-\beta(q-1)}} \{1 + i(1-q)(k+i\epsilon)\lambda^{\beta(q-1)}y\}^{\frac{1}{1-q}} dy = \\
& \frac{(q-1)[1-\beta(q-1)]a^{\frac{q-2}{q-1}}}{(2-q)[(1-q)i(k+i0)\lambda^{\beta}]^{\frac{1}{q-1}}} \times \\
& F\left(\frac{1}{q-1}, \frac{2-q}{(q-1)[1-\beta(q-1)]}, \frac{1}{q-1} + \frac{\beta(2-q)}{1-\beta(q-1)}; \right. \\
& \left. -\frac{1}{(1-q)i(k+i0)\lambda^{\beta(q-1)}a^{1-\beta(q-1)}}\right), \tag{3.6}
\end{aligned}$$

and

$$\begin{aligned}
& \lim_{\epsilon \rightarrow 0^+} \int_0^{a^{1-\beta(q-1)}} y^{\frac{\beta(2-q)}{\beta(q-1)-1}} \{1 + i(1-q)(k+i\epsilon)\lambda^{\beta(q-1)}y\}^{\frac{1}{1-q}} dy = \\
& \frac{[\beta(q-1)-1]a^{1-\beta}}{\beta-1} \times \\
& F\left(\frac{1}{q-1}, \frac{\beta-1}{\beta(q-1)-1}, \frac{\beta q-2}{\beta(q-1)-1}; \right. \\
& \left. (q-1)i(k+i0)\lambda^{\beta(q-1)}a^{1-\beta(q-1)}\right), \tag{3.7}
\end{aligned}$$

where  $F(a, b, c; z)$  is the hypergeometric function. Thus we obtain for  $F(k, q)$

$$\begin{aligned}
& F(k, q) = [H(q-1) - H(q-2)] \times \\
& \left\{ \left\{ H(q-1) - H\left[q - \left(1 + \frac{1}{\beta}\right)\right] \right\} \times \right. \\
& \left. \frac{(q-1)\lambda^{\beta}}{(2-q)[(1-q)i(k+i0)\lambda^{\beta}]^{\frac{1}{q-1}}} \times \right. \\
& \left. \left\{ a^{\frac{q-2}{q-1}} F\left(\frac{1}{q-1}, \frac{2-q}{(q-1)[1-\beta(q-1)]}, \frac{1}{q-1} + \frac{\beta(2-q)}{1-\beta(q-1)}; \right. \right. \right.
\end{aligned}$$

$$\begin{aligned}
& \left. \frac{1}{(q-1)i(k+i0)\lambda^{\beta(q-1)}a^{1-\beta(q-1)}} \right) - \\
& b^{\frac{q-2}{q-1}} F \left( \frac{1}{q-1}, \frac{2-q}{(q-1)[1-\beta(q-1)]}, \frac{1}{q-1} + \frac{\beta(2-q)}{1-\beta(q-1)}; \right. \\
& \left. \frac{1}{(q-1)i(k+i0)\lambda^{\beta(q-1)}b^{1-\beta(q-1)}} \right) \Big\} + \\
& \left\{ H \left[ q - \left( 1 + \frac{1}{\beta} \right) \right] - H(q-2) \right\} \frac{\lambda^\beta}{\beta-1} \times \\
& \left\{ a^{1-\beta} F \left( \frac{1}{q-1}, \frac{\beta-1}{\beta(q-1)-1}, \frac{\beta q-2}{\beta(q-1)-1}; \right. \right. \\
& \quad \left. \left. (q-1)i(k+i0)\lambda^{\beta(q-1)}a^{1-\beta(q-1)} \right) - \right. \\
& \quad \left. b^{1-\beta} F \left( \frac{1}{q-1}, \frac{\beta-1}{\beta(q-1)-1}, \frac{\beta q-2}{\beta(q-1)-1}; \right. \right. \\
& \quad \left. \left. (q-1)i(k+i0)\lambda^{\beta(q-1)}b^{1-\beta(q-1)} \right) \right\} \Big\}. \tag{3.8}
\end{aligned}$$

As we can see from (3.8),  $F(k, q)$  is dependent of  $a$  and  $b$ , and, as consequence, one-to-one as has been shown in Section 2.

However, and this is the crucial issue, if we **fix**  $q$  and select  $\beta = 1/(q-1)$

(3.8) simplifies and adopts the appearance

$$\begin{aligned}
F(k, q) &= \lambda^{\frac{1}{q-1}} \frac{q-1}{2-q} [H(q-1) - H(q-2)] \times \\
& \left[ a^{\frac{q-2}{q-1}} F \left( \frac{1}{q-1}, \frac{2-q}{q-1}, \frac{2-q}{q-1}; (q-1)i(k+i0)\lambda \right) - \right. \\
& \left. b^{\frac{q-2}{q-1}} F \left( \frac{1}{q-1}, \frac{2-q}{q-1}, \frac{2-q}{q-1}; (q-1)i(k+i0)\lambda \right) \right]. \tag{3.9}
\end{aligned}$$



With the help of the result given in [27] for

$$F(-a, b, b, -z) = (1 + z)^a,$$

we obtain for (3.9):

$$F(k, q) = \lambda^{\frac{1}{q-1}} \frac{q-1}{2-q} [H(q-1) - H(q-2)] \left( a^{\frac{q-2}{q-1}} - b^{\frac{q-2}{q-1}} \right) [1 + (1-q)ik\lambda]^{\frac{1}{1-q}}. \quad (3.10)$$

Using now the expression for  $\lambda$  of [22], i.e.,

$$\lambda = \left[ \left( \frac{q-1}{2-q} \right) \left( a^{\frac{q-2}{q-1}} - b^{\frac{q-2}{q-1}} \right) \right]^{1-q},$$

we have, finally,

$$F(k, q) = [H(q-1) - H(q-2)] [1 + (1-q)ik\lambda]^{\frac{1}{1-q}}, \quad (3.11)$$

which is the result given by Hilthorst in [22], demonstrating that  $F(k, q)$  is not one-to-one. As a conclusion we can say that for fixed  $q$  the  $q$ -Fourier transform is NOT one-to-one. On the contrary, as we have shown in section 2, when  $q$  is NOT fixed, the  $q$ -Fourier transform is indeed one-to-one.

## Conclusions

In the present communication we have discussed the NOT one-to-one nature of the  $q$ -Fourier transform  $F_q$ . We have shown that, if we eliminate the

requirement that  $q$  be fixed and let it “float” instead, such simple extension of the  $F_q$ –definition restores the desired one-to-one character.

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