Giant H II regions in NGC 7479 and NGC 6070

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ABSTRACT
We present new results from our search for giant H II regions in galaxies visible from the Southern hemisphere. In this work we study two galaxies: NGC 7479 and NGC 6070. Using high-resolution spectra, obtained with different instruments at Las Campanas Observatory, we are able to resolve the emission-line profile widths and determine the intrinsic velocity dispersion of the ionized gas. We detect profile widths corresponding to supersonic velocity dispersions in the six observed H II regions. We find that all of them show at least two distinct kinematical components: a relatively narrow feature (between 11 and 22 km s\(^{-1}\)) and a broader (between 31 and 77 km s\(^{-1}\)) component. Two of the regions show a complex narrow profile in all ion lines, which can be further split into two components with different radial velocities. Whereas the wing broadening of the overall profile can be fitted with a low-intensity broad component for almost all profiles, in one region it was better reproduced by two separate shell-like wings. We have analysed the impact that the presence of multiple components has on the location of the H II regions in the \(\log(L) - \log(\sigma)\) plane. Although the overall distribution confirms the presence of a regression, the precise location of the regions in the plane is strongly dependent on the components derived from the profile fitting.

Key words: H II regions – galaxies: individual: NGC 7479 – galaxies: individual: NGC 6070 – galaxies: starburst.

1 INTRODUCTION
The most extended and luminous objects which are observed on the discs of spirals, in irregulars and in starburst galaxies are called giant extragalactic H II regions (GH\,I\,Rs). GH\,I\,Rs are places of very active star formation and they are characterized by the large emission of ultraviolet photons from a large number of young and massive stars. Owing to the large number of ionizing photons most of the surrounding gas is ionized, and the recombination process provides a strong signature of recent or ongoing massive star formation.

The emission-line profile widths of the giant H II regions imply the existence of supersonic motions in the gas (Smith & Weedman 1970). Furthermore, Terlevich & Melnick (1981) found a correlation between the gas velocity dispersion (\(\sigma\)) and the total luminosity (\(L\)) emitted in the respective line. Hippelein (1986) and Arsenault & Roy (1988) confirmed the existence of such a regression, but there was no agreement on the values derived for its slope so no conclusion could be reached regarding the origins of this motion. Fuentes-Masip et al. (2000) and Bosch, Terlevich & Terlevich (2002) using different techniques on different samples agreed on a slope close to a value of 4, which seemed to favour the gravitational source of energy for the observed supersonic motion. Even if the nature of this behaviour is not fully understood, it still allows us to distinguish giant H II regions from an agglomeration of classical H II regions by means of high-resolution spectroscopy, GH\,I\,Rs provide the link between small-scale star-forming regions, such as Orion in our Galaxy and violent star formation processes taking place in starburst galaxies.

In Firpo, Bosch & Morrell (2005) we have confirmed the giant nature of three candidates to giant H II regions in the southern sky which were identified as very luminous H II regions by Feinstein (1997).

Here we continue our search and detailed analysis of GH\,I\,Rs in local Universe galaxies. We have selected from the Feinstein catalogue, a sample of the brightest H II regions from two spiral galaxies: NGC 7479 and NGC 6070.

NGC 7479 is a barred spiral galaxy, at a distance of 31.92 Mpc (Rozas et al. 1999), classified as SB(s)c by de Vaucouleurs et al. (1991), Devereux (1989) and Martinet & Friedli (1997) have classified it also as a starburst galaxy, whereas Keel (1983) and Ho, Filippenko & Sargent (1997) have identified it as a low-ionization nuclear emission-line region (LINER) and a Seyfert 1.9 galaxy, respectively. Many studies have been performed on this galaxy owing to the presence of a certain asymmetry in the spiral structure

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(Quillen et al. 1995) that leads several authors to suggest this might be owed to a recent merger event (Laine 1996; Laine & Heller 1999). This scenario can also account for perturbations in the velocity field (Laine & Gottesman 1998) and the latest minor merger model indicates that the remnant may be situated within the bar (Laine 2001) as no remnant of the merger could be identified in optical images (Saraiva & Benedict 2003).

NGC 6070, at a distance of 29.8 Mpc (García-Gómez, Athanassoula & Barbé 2002), is classified as SA(s)cd by de Vaucouleurs et al. (1991) although Gresbol, Patsis & Pompei (2004) state that a bar with two main arms can be identified. Pohlen & Trujillo (2006) further studied the bar structure classifying it as Type II.o,CT according to Erwin, Pohlen & Beckman (2008). The galaxy shows a normal arm distribution, with two arms that break into several on each side of the central region (see Fig. 2). The disc dynamics was studied by Márquez et al. (2002), and García-Gómez et al. (2002) studied the distribution of H α regions in the galaxy.

In order to verify the nature of our giant H α region candidates, we obtained high-resolution spectra to measure the emission-line profile widths and to estimate if the velocity dispersion is indeed supersonic. To determine the optimum number of Gaussian components to fit each line profile in order to minimize the uncertainty in the velocity dispersion of the ionized gas, we have considered information from other studies available in the literature. Previous papers have dealt with the fact that a single Gaussian profile might not be realistic enough when fitting the observed emission lines. Several studies favour the presence of an omnipresent broad component which explains the integral profile wings. They include Melnick, Tenorio-Tagle & Terlevich (1999) who identified a broad component in the combination lines in the central region of the 30 Doradus nebula, Díaz et al. (1987) and Terlevich et al. (1996) in the M33 giant H α region NGC 604, Mendez & Esteban (1997) in four Wolf-Rayet galaxies and Homeier & Gallagher (1999) in the starburst galaxy NGC 7673. More recently, Hägele et al. (2007, 2009, 2010) and Hägele (2008) found that the best Gaussian fits in circumnuclear star-forming regions involved the presence of broad and narrow components for the emission lines of the ionized gas.

A variety of line broadening mechanisms have been proposed to interpret the existence of the broad supersymmetric component measured in the emission-line profile of GH αRs. Dynamics of virialized systems (Tenorio-Tagle, Muñoz-Tuñón & Cox 1993), superposition of multiple gas bubbles in expansion (Chu & Kennicutt 1994) or turbulence of the same interstellar gas (Medina Tanco et al. 1997). More recently, Westmoquette et al. (2007a,b) concluded that the narrow component represents the general disturbed ionized interstellar medium (ISM), arising through a convolution of the stirring effects of the starburst and gravitational virial motions. On the other hand, the broad component results from the highly turbulent velocity field associated with the interaction of the hot phase of the ISM with cooler gas knots, setting up turbulent mixing layers.

Other studies support the existence of shell-like profiles on the blue and red wings of the main feature as found by Chu & Kennicutt (1994), who state that the integrated line profiles of the entire 30 Dor region may be fitted with a broad Gaussian with low-intensity wings. Relaño & Beckman (2005) claim that a large fraction of the H α regions in three barred spiral galaxies shows high-velocity and low-intensity features in its integrated line profiles and Rozas et al. (2006) find that the great majority of their sampled H α regions show that the brightest principal component is accompanied by low-intensity components symmetrically shifted in velocity.

In this paper we present echelle data obtained at Las Campanas Observatory (LCO) of six GH αRs candidates in the two mentioned galaxies. We have determined the velocity dispersion and we have evaluated the possible presence of a broad component or two symmetric low-intensity components in the fit with the observed emission-line profile widths. In Section 2 we present the observations and the data reduction. Section 3 details the analysis performed on the emission-line profiles and discusses each conspicuous case individually. The summary and conclusions of this work are included in Section 4.

2 OBSERVATIONS AND REDUCTIONS

According to canonical denomination for GH αRs we have labelled our candidates following their brightness ranking from Feinstein catalogue, NGC 7479 I is therefore the brightest region in NGC 7479, NGC 7479 II is the second brightest and so on. Observed regions from NGC 7479 and NGC 6070 are shown in Figs 1 and 2, respectively. High-resolution spectra were obtained with an echelle spectrograph at the 100-inch du Pont Telescope, LCO, between 2006 July 19 and 22. The spectral range covered by the observations was from 3800 to 9500 Å. Observing conditions were good with 1 arcsec seeing and photometric sky. A 2 × 2 binning was applied to the CCD in order to minimize readout contribution to the final spectrum noise. The spectral resolution achieved in our du Pont Echelle data (obtained with 1 arcsec effective slit width) is $R \approx 25,000, \Delta \lambda = 0.25 \lambda$ at $\lambda 6000 \AA$, as measured from the full width at half-maximum (FWHM) of the Th–Ar comparison lines. This translates in a resolution of $\sim 12 \text{ km s}^{-1}$.

The standard stars, HR 7950 and HR 4468, and the CALSPEC spectrophotometric standard star, Feige 110 (Bohlin, Dickinson & Calzetti 2001), were also observed for flux calibration purposes. In addition, Th–Ar comparison spectra, milky flats (sky flats obtained with a diffuser, during the afternoon) and bias frames were taken every night. The exposure time for standard stars was 5 s for the bright HR stars and 1200 s for Feige 110.
We also compared the red end of the wavelength-calibrated spectra with the night-sky spectrum by Osterbrock et al. (1996). This turned out to be a very reliable confirmation of the goodness of the wavelength solution, and we were able to check that differences between our wavelengths and the sky line wavelengths were below 0.05 Å.

We performed flux calibrations using the observed spectrophotometric standard stars as described in Firpo et al. (2005). The observed standard stars had their fluxes tabulated every 16 A, and the amount of defined intervals within an echelle order ranged from four to 12, depending on the quality of the spectrum.

For the Magellan spectra, and in the case of the NGC 7479 I region du Pont data, we were able to obtain CALSPEC spectrophotometric standard stars (Bohlin et al., 2001), whose fluxes were tabulated every 2 Å. These stars are ideal for calibrating high-resolution echelle spectra, although their relatively low brightness (Feige 110, V = 11.83 and BD+28D4211, V = 10.51) makes them time-consuming targets for 2-m class telescopes.

### 3 ANALYSIS OF LINE PROFILES

Making use of the known redshifts z = 0.006685 for NGC 6070 (Springob et al. 2005) and z = 0.007942 for NGC 7479 (Haynes et al. 1998) we identify the hydrogen recombination lines, such as Hα and Hβ and collisionally excited lines, such as [N II] λ6548, 6584 Å, [S II] λ6717, 6731 Å, present in the spectra. These strong lines are used to analyse the structure of the profiles as they allow us to verify the existence of more than one component. The adopted laboratory wavelengths are taken from work by García-Rojas et al. (2005).

To determine the radial velocities and the velocity dispersions of the ionized gas, we measure the central wavelength and width of several emission lines. The radial velocity of each emission line is obtained from its central wavelength determined from the Gaussian profile fitting, and their errors are given by the fitting error provided by the resampling done within the xGAUSS task of IRAF and taking into account the rms of the wavelength calibration. The true velocity dispersion (σ) of each emission line is calculated as described in Firpo et al. (2005). The width of the observed profile (σo) is affected by the contribution of thermal random motions (σt) and the instrumental profile (σi). Then, we can obtain the true velocity dispersion as σ^2 = σ^2_0 - σ^2_t - σ^2_i. As in Firpo et al. (2005), a typical kinetic temperature T = 10^4 K is assumed, and the instrumental profile (σi) is very well approximated by a single Gaussian function. The velocity dispersion errors have been calculated using the observational errors in the observed profile and assuming negligible errors in σt and σi. Owing to the generally high metallicity of these kinds of objects as a class, the temperatures are probably lower than 10^5 K (see Díaz et al., 2007, and references therein). The low excitation of these objects estimated from their spectra makes any temperature sensitive line too weak to be observed. In fact, in most cases, the [O III] λ5007 Å line, which is typically 100 times more intense than the auroral [O III] λ4363 Å one, can barely be seen. In this way, we have also calculated the velocity dispersions assuming a kinetic temperature of 5000 K and we do not find great differences with the previous calculations derived using 10^4 K. These differences are between 0.02 and 3 km s^-1 depending on the measured velocity dispersion since σi is constant for a given temperature and subtracted in quadrature to σo. Obviously, the estimated σ are always higher for the lower assumed kinetic temperatures. Then, for consistency with Firpo et al. (2005) we use a kinetic temperature of 10^4 K. Therefore, this effect owing to the kinetic temperature

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1Image Reduction and Analysis Facility, distributed by NOAO, operated by AURA, Inc., under agreement with NSF.
uncertainty, would produce an underestimation in the derived true velocity dispersions.

When fitting single Gaussian profiles to the emission lines we notice that, although it works well for most of the observed profiles, there is a residual present in the wings of several lines. This has been observed in GH\Rs before. Some authors (Melnick et al. 1999; Häggle et al. 2007, 2009, 2010, among others) have shown that the profile wings are well fitted by single (broad) Gaussian of supersonic widths, whereas other authors (Chu & Kennicutt 1994; Relaño & Beckman 2005; Rozas et al. 2006) proposed that the profiles appear to be dominated by a few, well-resolved, expanding shells. Wherever possible, we have considered and applied both suggestions.

In order to fit multiple components to the observed profiles we make use of the \textsc{negaussfit} task within \textsc{iraf}. \textsc{negaussfit} performs iterative fitting of multiple Gaussian profiles to spectral lines. The task allows us to select individual parameters to be fitted and delivers an estimate for the uncertainty of derived parameters. The task needs initial guesses for the function coefficients, and these are specified through previously generated tables.

In those cases when more than one component is evident, we follow an iterative procedure in which we only allow the task to fit a limited subset of parameters at one time. This is done to constrain the universe of possible solutions to the fit by making use of the different information available from each emission line. Basically, after obtaining an initial guess for the second component parameters, we fit the amplitude, centre and width of each component at one time, leaving the parameters of the other one untouched. Once the central wavelength of all components is relatively well known, we set them as fixed parameters and allow the task to perform a final fit of amplitude and width of all components at the same time.

The profile fittings show the presence of a residual in the emission-line wings for all cases, becoming more evident in the strong Hα emission line. Therefore, we evaluate the presence of either a broad component or shell features. For the former assumption, we introduce a broad component with an initial width which is three times the width of the narrow component(s). This guess comes from an overall view of the values found by Häggle et al. (2007) when performing a similar analysis. We fit the peak values for the narrow components and then their profile widths. The broad component parameters are then fitted before we start the iterative procedure again. In the latter scenario, a similar iterative procedure takes place, but in this case we add a blue and a red narrow component to account for the additional emission. The validity of the profile multiplicity and broadening is checked over the different emission lines available for each region, although the most reliable profile analysis comes from the strongest lines, such as Hα.

Out of our sample of six GH\Rs we find that all of them show evidence of wing broadening. For these regions we attempt to fit an overall broad component or two narrow components symmetrically shifted in velocity with respect to the intense component. In the following subsections we will discuss our findings grouping them by the outcome profiles.

### 3.1 NGC 7497 I

Rozas et al. (1999) undertook a complete study to date of the H\β regions in NGC 7497 galaxy, determining positions, angular sizes and absolute fluxes of over 1000 H\β regions and they constructed the luminosity function for the regions over the whole galaxy. We have cross-correlated these results with our sample and find that the three NGC 7497 regions selected in our work can be identified in the Rozas catalogue: NGC 7497 I as RZH99 = 1, NGC 7497 II

### Table 2. Results of Gaussian profiles fitting to the observed emission lines in NGC 7497. Each emission line is identified by its ion laboratory wavelength and ion name in columns 1 and 2. According to the different fits performed on each line, column 3 identifies each 'narrow component' (A and B, where applicable), a broad component and/or blue and red wings (b wing and r wing, respectively). Radial velocities ($V_r$) and intrinsic velocity dispersions ($\sigma_{int}$) together with their respective errors are expressed in km s$^{-1}$. The intrinsic velocity dispersions are corrected for the instrumental and thermal widths. EMs are included and shown as a percentage of the component flux relative to the total EM of the region.

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as RZH99 = 7 and NGC 7479 III as RZH99 = 74 (RZH99 is the number identification used by Rozas et al. 1999).

We identify and fit the Gaussian profiles to the Hα, [N ii] λ6584 Å and [S ii] λλ6717, 6731 Å lines in NGC 7479 I region with splot routines. The Gaussian fits reveal the presence of two distinctly separated kinematic components (labelled A and B) in all analysed profiles. Component A shows a much broader profile than component B, and both components show little spread in their individual radial velocities among different emission lines present in the spectrum. The reliability of these values is confirmed when we improve the profile fitting using ngaussfit which in turn yields values for the profile width of each component.

Even considering the presence of two distinct components with different radial velocities, the overall fit cannot account for a residual emission in the profile wings. Following the procedures outlined in the previous section, we are able to fit a broad component, with a velocity dispersion of about 75 km s\(^{-1}\) from the Hα emission line.

![Figure 3](image)

**Figure 3.** nGauss fits with three Gaussian components in the NGC 7479 I emission-line profiles, two narrow principal components and a broad one. In order: Hα, [N ii] λ6584 Å and [S ii] λ6731 Å.

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![Figure 4](image)

**Figure 4.** 2D Hα velocity image of NGC 7479 I region with the [N ii] λ6584 Å (red) contours overlapped. Three sections are numbered in the 2D image.

**Table 3.** Gaussian fitting results for individual sections of the Hα emission-line profile in NGC 7479 I. Kinematical parameters [radial velocity (\(V_r\)) with errors (km s\(^{-1}\))] the intrinsic velocity dispersions (\(\sigma_{\text{int}}\)) corrected for the instrumental and thermal widths (km s\(^{-1}\)) and EM of each component as a fraction of the total EM of the sampled section (per cent) data are shown.
Figure 5. NGAUSS fits with three (two) Gaussian components in the NGC 7479 I sections, two (one) narrow principal components and a broad one. In order: Hα (sections 2, 4, 5 and 7), [N II] λ6584 Å (sections 4 and 5).
2540 km s\(^{-1}\) labelled A and B, respectively. Each of these relatively narrow features has a broader component associated with it. Although the narrow component can still be deblended in the central region, we are not able to disentangle the broad ones, which results in a single broad profile with a velocity similar to the average between components A and B. Regarding the velocity dispersion of each component, we have listed them in Table 3 as a reference, but they should not be taken as a definite value as the signal-to-noise ratio is variable and somewhat low for the slit edges.

If we compare these results with the radial velocity map derived for the ionized gas in NGC 7479 (Laine 1996) we find, from visual inspection, that the expected velocity at the position of NGC 7479 I ranges between 2520 and 2535 km s\(^{-1}\) close to an average value of our components A and B. A similar comparison, using the H\(_I\) map by Laine & Gottesman (1998) which has a better spatial and spectral resolution, yields a value between 2530 and 2540 km s\(^{-1}\) much closer to the value measured for component B. Although the latter map is sampling neutral rather than ionized hydrogen, the former might be affected by the complex structure we find within the region. At this stage we can only suggest that component A shows an odd kinematic behaviour and that a detailed study of NGC 7479 I and its surroundings is needed to clarify this issue.

3.2 NGC 7479 II

A similar procedure is used to analyse the profiles of the emission lines of NGC 7479 II. In this case, two components are identified with different radial velocities, but as the overall signal obtained

![Figure 6. NGAUSS fits with four Gaussian components in the NGC 7479 II emission-line profiles, two narrow principal components and two low-intensity components. In order: H\(_\alpha\) and [N \text{II}] 6584 Å.](image)

![Figure 7. NGAUSS fits with two Gaussian components in the NGC 7479 III emission-line profiles. In order from left to right and from top to bottom: H\(_\alpha\), [N \text{II}] 6584 Å, [S \text{II}] 6717 Å and [S \text{II}] 6731 Å.](image)
Figure 8. *ngauss* fits with two Gaussian components in the NGC6070I emission-line profiles. du Pont data on the left-hand side and MIKE data on the right-hand side. In order from top to bottom panels: [O iii] λ5007 Å, Hα, [N ii] λ6584 Å and [S ii] λ6717 Å. *ngauss* fits with two Gaussian components in the NGC 6070I emission line profiles. du Pont data on the left-hand side and MIKE data on the right-hand side for [S ii] λ6731 Å line.
for this region is lower than for NGC 7479 I it is not possible to disentangle them with 2D spectra. These two narrow components are found in the Gaussian profiles of the Hz, Hβ, [N II] λλ6548, 6584 Å and [S II] λλ6717, 6731 Å lines. In this region, however, the wings of the brightest component profile could not be fitted with one single broad component. On the other hand, we are able to fit it with two wings only for the two lines with the strongest signal, namely Hα and [N II] 6584 Å (see Fig. 6). The final fitting parameters are shown in Table 2. The overall Hα flux, uncorrected for reddening, is found to be 1.06 × 10^{-13} erg s^{-1} cm^{-2}. The emission lines with a low signal-to-noise ratio for which the fits do not provide accurate results for the broad and the two wings components are not listed in the table.

We compare the velocity separation between the redshifted (and blueshifted) components with the average of the central features and the Hα luminosity of the region. The velocity separation of the red- and blueshifted components in this work is in agreement with the regression found by Relano & Beckman (2005) for their HII regions sample.

We also analyse the radial velocity behaviour within the galaxy, and observe a similar situation to that of NGC 7479 I. There is no difference, within the errors, when we compare our average value derived from the Hα emission line with the Hα velocities field map (Laine 1996). If, however, the comparison is done against the HII rotation velocity map (Laine & Gottesman 1998), knot B seems to follow galactic rotation whereas knot A does not.

### 3.3 NGC 7479 III

The observed spectrum of this region shows, although weak, the [O III] λλ4959, 5007 Å lines. This is probably owed to a difference in metallicity, NGC 7479 III being slightly less metallic than the brightest two regions discussed above, or a difference in the ionization structure. The latter might be due to a relatively small difference in the hardness of the ionization radiation field (and, what is equivalent, in the ionizing stellar population) responsible for the excitation observed in the spectra or a difference in the nebular geometry (see a detailed discussion in Hägele et al. 2006, 2008; Díaz et al. 2007; Pérez-Montero et al. 2010, and references therein). Unfortunately, as we said above, since no temperature sensitive lines can be observed, we are not able to derive their metallicities, and distinguish the origin of this small difference in the excitation present in the spectra of these regions of NGC 7479.

We identify and fit the Gaussian profiles to the [O III] λ5007 Å, Hα, [N II] λ6584 Å and [S II] λλ6717, 6731 Å lines in NGC 7479 III region with nGaussfit. In this region we have found one supersonic component, with no evidence of multiple contributions but it is possible to detect flux excess on the wings of the emission lines. We therefore fit a broad Gaussian component together with the narrow component (see Fig. 7). This procedure works well for almost all detected lines. In the case of [O III] λ5007 Å, however, there are no prominent residuals when we fit only one Gaussian component which has a dispersion of about 20 km s^{-1} and we are therefore not able to fit a broad Gaussian component. For the other listed lines the derived radial velocities, the velocity dispersions (using σ^2 = 5.4 km s^{-1} and T ≈ 10^4 K) and errors are listed in Table 2. The overall Hα flux, uncorrected for reddening, is found to be 5.43 × 10^{-13} erg s^{-1} cm^{-2}. It is worth mentioning that there seems to be some systematic offset between the derived profile centres of the narrow and broad profiles which is, in average, about 10 km s^{-1}. This happens for all lines, although it becomes more evident for [S II] λ6731 Å. An undiscovered blueshifted narrow component could account for the observed behaviour, although better signal-to-noise ratio data are needed to confirm this hypothesis.

The radial velocity average derived for the single component is in agreement with the expected radial velocity at the position of NGC 7479 III, showing only a slight offset of (10 km s^{-1}), well within the uncertainties.

### 3.4 NGC 6070 I, II and IV

All regions observed in NGC 6070 show one narrow component together with the underlying broad component. We identify and fit the Gaussian profiles to the Hβ, [O III] λ5007 Å, Hα, [N II] λ6584 Å and [S II] λλ6717, 6731 Å lines in almost all regions (see Figs 8–10). Table 4 shows the derived kinematical parameters (using σ^2 = 5.3, 5.6 and 5.3 km s^{-1} for regions I, II and IV, respectively; T is set to ≳10^4 K for all regions).

For these three regions, we are also able to compare the du Pont data with higher resolution spectra data obtained with the double echelle spectrograph MIKE in 2004 July. The agreement in the kinematical parameters can be seen in Table 4, which shows the reliability of the data here presented. The overall Hα flux, uncorrected for reddening, for du Pont data are 1.68 × 10^{-13}, 6.26 × 10^{-14} and 2.15 × 10^{-13} erg s^{-1} cm^{-2} for regions I, II and IV, respectively, and for MIKE data are 1.74 × 10^{-13}, 1.29 × 10^{-13} and 1.88 × 10^{-13} erg s^{-1} cm^{-2}. Only a slight bias towards lower velocity dispersion can be seen for NGC 6070 I, which can be probably attributed to differences in the centring of the slit during acquisition of the target. There seems to be some systematic offset between
Figure 9. \textit{ngauss} fits with two Gaussian components in the NGC 6070 II emission-line profiles. du Pont data on the left-hand side and MIKE data on the right-hand side. In order from top to bottom panels: H$\alpha$, H$\alpha$, [N II] $\lambda$6584 Å and [S II] $\lambda$6731 Å.
Figure 10. *sgauss* fits with two Gaussian components in the NGC 6070 IV emission-line profiles, du Pont data on the left-hand side and MIKE data on the right-hand side. In order from top to bottom panels: [O III] λ5007 Å, Hα, [N II] λ6584 Å and [S II] λ6717 Å.
Table 4. Results of Gaussian profiles fitting to the observed emission lines in NGC 6070. Each emission line is identified by its ion laboratory wavelength and ion name in columns 1 and 2. According to the different fits performed on each line, column 3 identifies each 'narrow component' (A and B, where applicable), a broad component and/or blue and red wings (b wing and r wing, respectively). Radial velocities ($V_r$) and intrinsic velocity dispersions ($\sigma_{int}$) together with their respective errors are expressed in km s$^{-1}$. The intrinsic velocity dispersions are corrected for the instrumental and thermal widths. EMs are included and shown as a percentage of the component flux relative to the total EM of the region. Upper panel: du Pont data; lower panel: MIKE data.

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</tbody>
</table>
| 4860 | H\(\beta\) | Narrow | ... | ... | ... | ... | 1883.4 | 1.1 | 14.4 | 2.3 | 59 | ... | ... | ... | ... | ... | ...
| 5007 | [O iii] | Narrow | 1871.1 | 0.8 | 18.4 | 1.1 | 59 | ... | ... | ... | ... | ... | 1859.3 | 0.7 | 15.0 | 1.0 | 66
| 6563 | H\(\alpha\) | Narrow | 1868.3 | 0.6 | 17.2 | 0.8 | 69 | 1886.4 | 0.6 | 13.6 | 0.8 | 60 | 1858.0 | 0.6 | 12.2 | 0.8 | 62
| 6584 | [N ii] | Narrow | 1869.7 | 0.6 | 16.6 | 1.1 | 47 | 1885.6 | 0.9 | 17.2 | 0.9 | 55 | 1863.4 | 0.8 | 13.5 | 0.9 | 54
| 6717 | [S ii] | Narrow | 1865.2 | 1.7 | 21.2 | 3.2 | 55 | ... | ... | ... | ... | ... | 1857.0 | 1.4 | 16.8 | 2.6 | 49
| 6731 | [S ii] | Narrow | 1864.7 | 1.0 | 14.3 | 1.3 | 33 | 1883.5 | 1.3 | 23.1 | 1.9 | 79 | ... | ... | ... | ... | ...
| MIKE data | | | | | | | | | | | | | | | | | | |
| 4860 | H\(\beta\) | Narrow | ... | ... | ... | ... | 1890.9 | 0.3 | 20.0 | 0.4 | 61 | ... | ... | ... | ... | ... | ...
| 5007 | [O iii] | Narrow | 1868.4 | 0.3 | 13.6 | 0.4 | 49 | ... | ... | ... | ... | ... | 1857.3 | 0.3 | 16.7 | 0.4 | 64
| 6563 | H\(\alpha\) | Narrow | 1866.9 | 0.2 | 13.4 | 0.3 | 62 | 1890.4 | 0.2 | 18.4 | 0.3 | 62 | 1856.8 | 0.2 | 11.9 | 0.3 | 58
| 6584 | [N ii] | Narrow | 1870.2 | 0.2 | 12.3 | 0.4 | 39 | 1892.6 | 0.2 | 21.3 | 0.3 | 68 | 1861.8 | 0.3 | 10.3 | 0.6 | 34
| 6717 | [S ii] | Narrow | 1867.3 | 0.4 | 29.5 | 0.6 | 61 | 1902.7 | 1.1 | 62.5 | 1.1 | 32 | 1855.6 | 0.6 | 27.0 | 0.9 | 66
| 6731 | [S ii] | Narrow | 1868.6 | 0.5 | 11.3 | 1.0 | 26 | 1889.1 | 0.3 | 23.0 | 0.5 | 70 | ... | ... | ... | ... | ...
| | Broad | 1866.4 | 0.7 | 27.1 | 0.8 | 74 | 1896.7 | 2.5 | 49.8 | 4.4 | 30 | ... | ... | ... | ... | ...

the derived profile centres of the narrow and broad profiles. These values are, in average, about $-25\text{ km s}^{-1}$ for du Pont data and $-10\text{ km s}^{-1}$ for MIKE data in the case of NGC 6070 II, and about 10 and 5 km s$^{-1}$ for NGC 6070 IV. Regarding radial velocities of the three observed regions in NGC 6070, there is no detailed velocity map in the literature. We are able to compare them with rotation curves derived from long slit observations made by Márquez et al. (2002). Although the orientation of the slits does not match exactly the location of all our regions, the velocities are in fair agreement considering the angular distance of the H\(\Pi\) regions to the centre of NGC 6070.

3.5 Relation between H\(\alpha\) luminosities and velocity dispersion

We have also analysed the impact that the presence of multiple components has on the location of H\(\Pi\) regions in the log\((L)-\log(\sigma)\) plane. To evaluate this we have plotted in Fig. 11 the luminosities and velocity dispersions derived directly from our spectrophotometric data. Luminosities were derived from the fluxes measured directly from the component fitting to our echelle spectra (uncorrected by reddening), and using distances as published by Rojas et al. (1999) for NGC 7479 and by García-Gómez et al. (2002) for NGC 6070. We are aware that slit spectra might not provide the most realistic sample of the overall flux of the ionized nebula, but the accuracy of acquisition and guiding is not good enough as to allow us to perform a proper identification with published luminosity values in fields relatively ‘crowded’ with H\(\Pi\) regions by Feinstein (1997) and Rojas et al. (2006). Velocity dispersion values and its errors are taken directly from Tables 2 and 4. The luminosity errors have been estimated taking into account the errors in the amplitude ($A$) and the FWHM in the component fitting ($F = 1.0645 \times A \times \text{FWHM}$), and the distance errors. As a reference value we have plotted a few giant H\(\Pi\) regions from Bosch et al. (2002) together with their linear fit to their ‘young’ H\(\Pi\) regions.

Inspection of Fig. 11 reveals that the distribution of points in the log\((L)-\log(\sigma)\) plane are strongly dependent on the components derived from the profile fitting. Individual components, labelled A and B (when applicable) as in their respective tables, and global values are those derived from a single Gaussian fitting to the line profile. As expected, individual components have smaller fluxes and velocity dispersions than the global profile and points are therefore shifted in the diagram. The incidence of this analysis has, however, a different outcome for regions in NGC 6070 and NGC 7479. In NGC 6070 the values derived from the global profile seem to lie, within observational errors, in the expected location if they follow the relation expected for virialized systems. When the individual components are plotted in the same diagram, only NGC 6070 II lies close to the regression, and NGC 6070 I and IV seem to be too

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lines. These images show that the two narrow components are
spatially resolved and the kinematical information helps to split the
knot in the velocity–distance plane. There is also evidence of the
presence of a broad component in all the narrower extracted sec-
tions, at least in the Hα emission line. For NGC 7479 II region we
can also fit two relatively narrow Gaussian components. In both
regions the radial velocity average value of our components A and
B derived from the Hα emission line matches the global velocity in
the galaxy when we project our value over the isovelocity contours
in Hα velocity field map. A similar comparison, using the Hβ map
yields a value much closer to the value measured for component B.
Then, we could only suggest that component A shows an odd kine-
matic behaviour. New observations with better spatial and spectral
resolution are needed to clarify this point.

In the rest of the studied Hα regions, we find one narrow Gaussian
component together with an underlying broad component, with
no evidence of multiple narrow contributions. The radial velocity
average derived for the narrower component corresponds to the
one expected from galaxy rotation curves, within the observational
ersors, in all of these cases.

In most cases, the velocity dispersions of the broad components
for these Hα regions derived by optimal Gaussian fit are in agree-
ment with Hägele et al. (2009, 2010) for circumnuclear star-forming
regions with σv in Hβ line ~34 to 65 km s⁻¹, and those derived by
Chu & Kennicutt (1994) and Melnick et al. (1999) for 30 Doradus
nebula with σv in Hα line ~45 km s⁻¹.

The estimated offsets between the narrow and broad components
in NGC 7479 III, NGC 6070 II and NGC 6070 IV are in complete
agreement with the values found in NGC 2903 and NGC 3310 by
Hägele et al. (2009, 2010), respectively, which are between ~25
and 35 km s⁻¹.

The distribution of the regions in the log(L)–log(σ) plane are
strongly dependent on the components derived from the profile
fitting. Individual components have smaller fluxes and velocity dis-
persions than the global profile and points are therefore shifted in
the diagram. In NGC 6070 the values derived from the global pro-
file seem to lie, within observational errors, in the expected location
if they follow the relation expected for virialized systems. When
the individual components are plotted in the same diagram, only
NGC 6070 II lies close to the regression, and NGC 6070 I and IV
seem to be too luminous for their measured velocity dispersion.
On the other hand, two out of three regions in NGC 7479 lie well away
this regression if parameters are derived from a single Gaussian, but
four out of five individual components show a very good agreement
when plotted in the same diagram. Definitely, the presence of more
than one Gaussian component rules the final position of the Hα
regions in the log(L)–log(σ) plane.

4 SUMMARY AND CONCLUSIONS

From new high-resolution spectra of the Hα regions NGC 7479 I,
NGC 7479 II, NGC 7479 III, NGC 6070 I, NGC 6070 II and
NGC 6070 IV obtained at the 100-inch du Pont Telescope, LCO,
we have confirmed the giant nature of all these regions. We have
also found that all of them show evidence of wing broadening evi-
dent mainly in the Hα line and confirmed in other emission lines.
In five of them we are able to fit a broad component which explains
the integral profile wings. Only in NGC 7479 II can we fit the integrated
profile wings with two narrow components symmetrically shifted in
velocity with respect to the average of the two intense components.

NGC 7479 I and NGC 7479 II reveal the presence of two sepa-
rated kinematic components with relatively narrow profiles in all
analysed emission lines. For NGC 7479 I region we can also gen-
erate 2D velocity images in Hα, [N II] λ6584 A and [S II] λ6717 A

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Space Administration.

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